

# Investigating magnetised, radiatively driven plasmas with a university scale pulsed-power generator

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**Imperial College (CIFS, Computational):** A. Crilly, J. Chittenden, S. Rose

**University of Nevada, Reno:** R. C. Mancini

**Imperial College  
London**

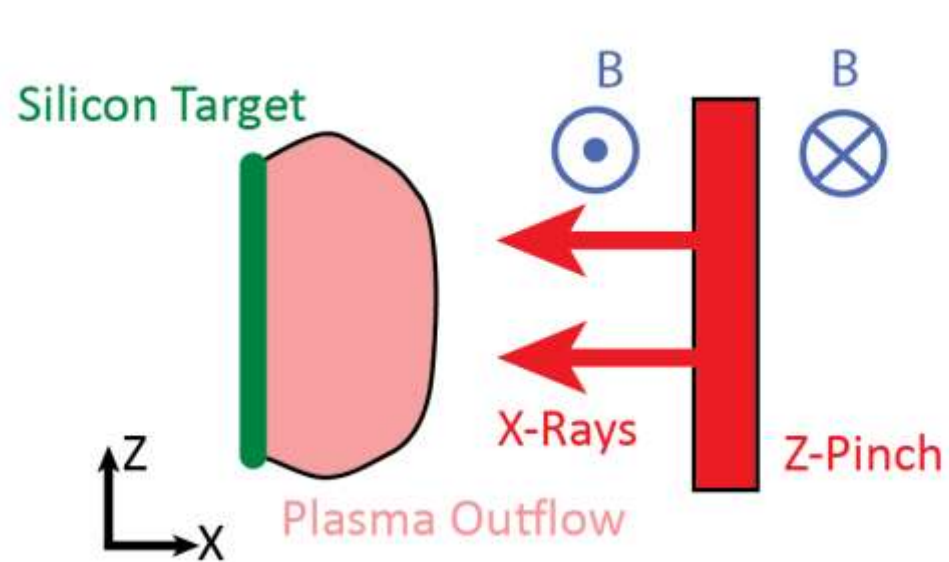


**MAGPIE**

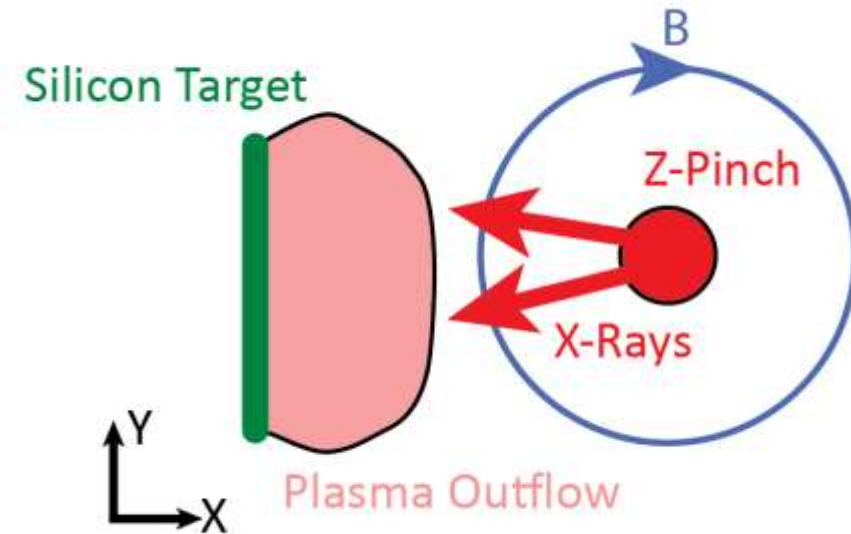


University of Nevada, Reno

# Overview of experimental setup



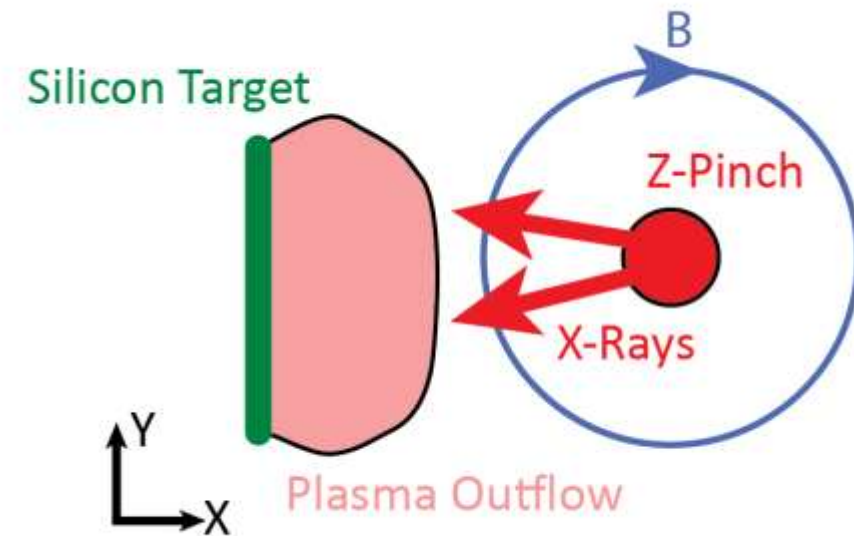
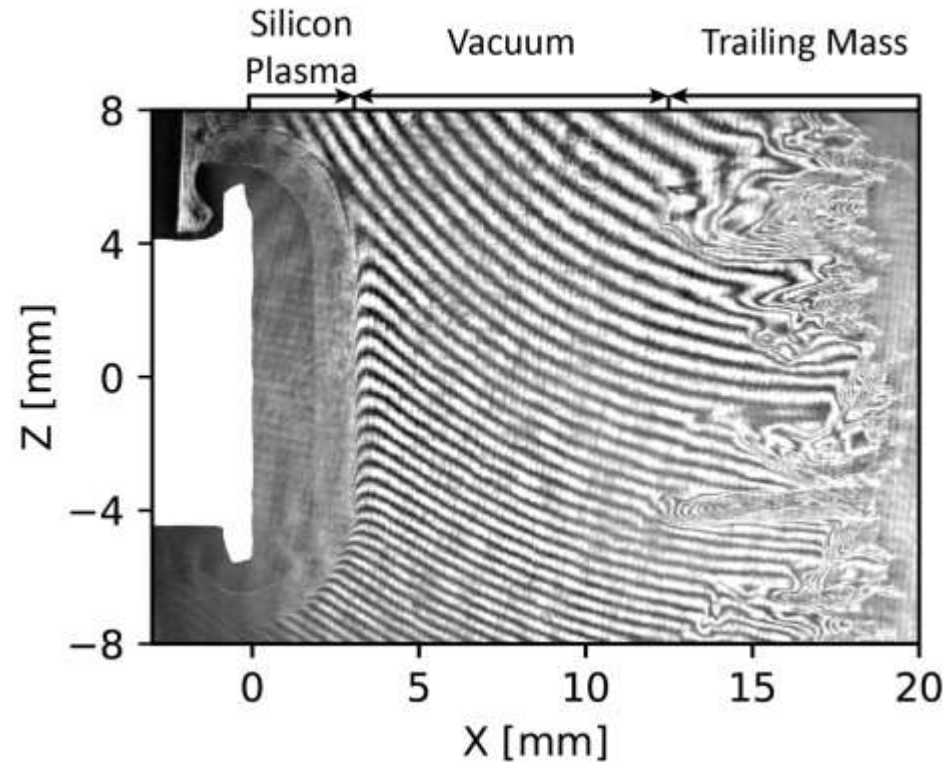
**Side-On** (X-Z plane) view of the experiment



**End-On** (X-Y plane) view of the experiment

- X-Rays from aluminium wire array Z-Pinch
- Ablated silicon plasma expands into  $\sim 10$  T magnetic field
- Experiments driven by MAGPIE (1.4 MA, 240 ns)
- Target positioned 1.5 – 4 cm from pinch

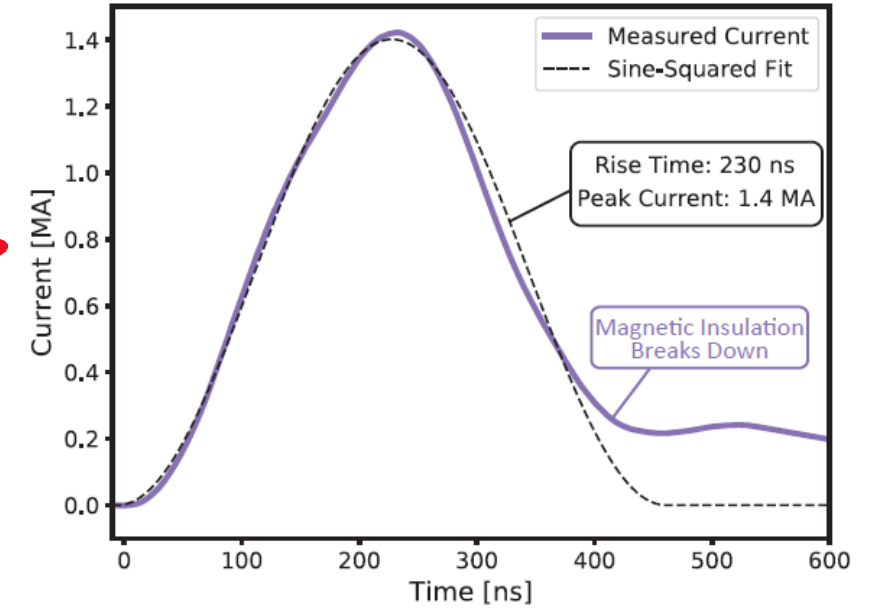
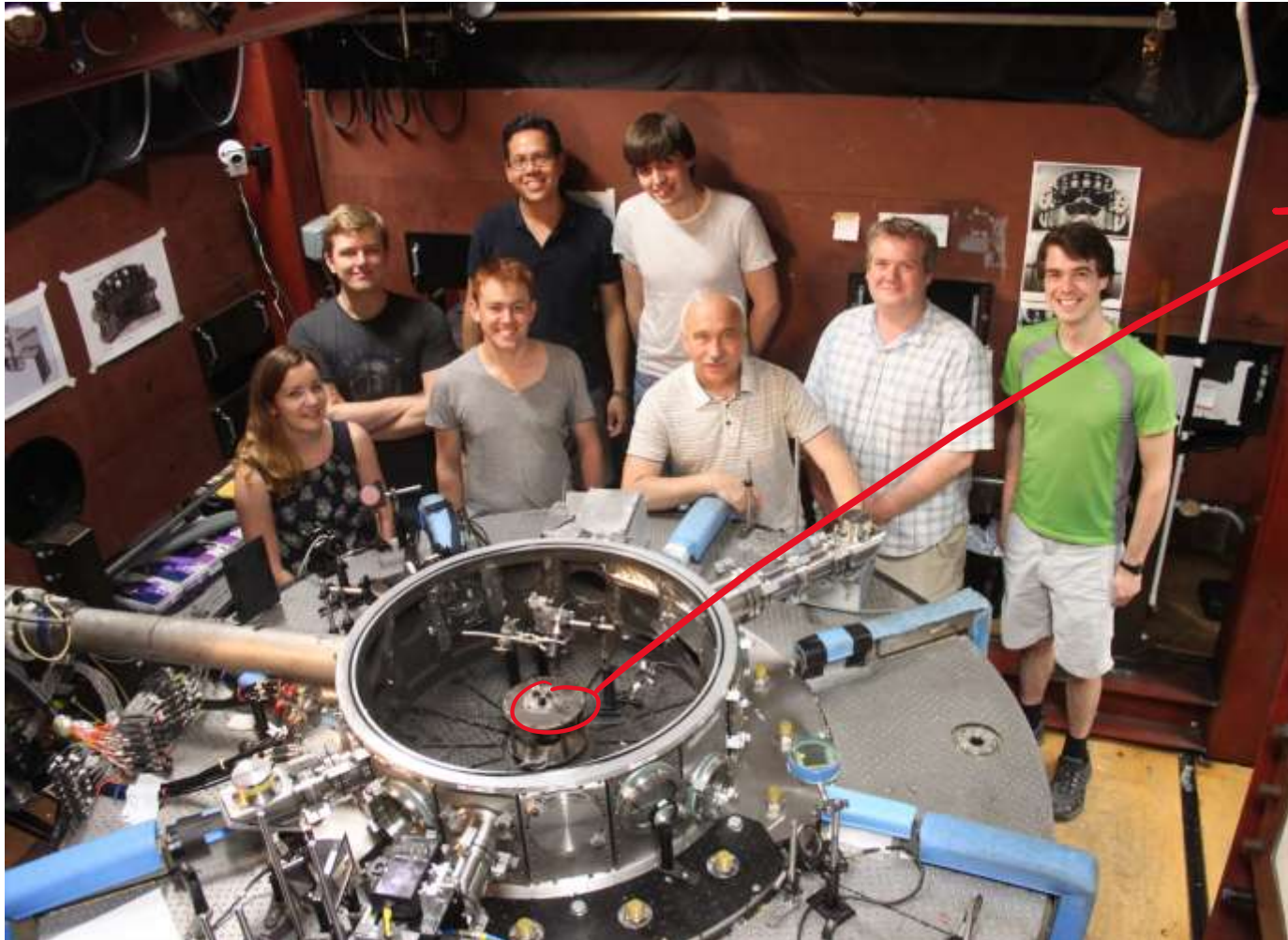
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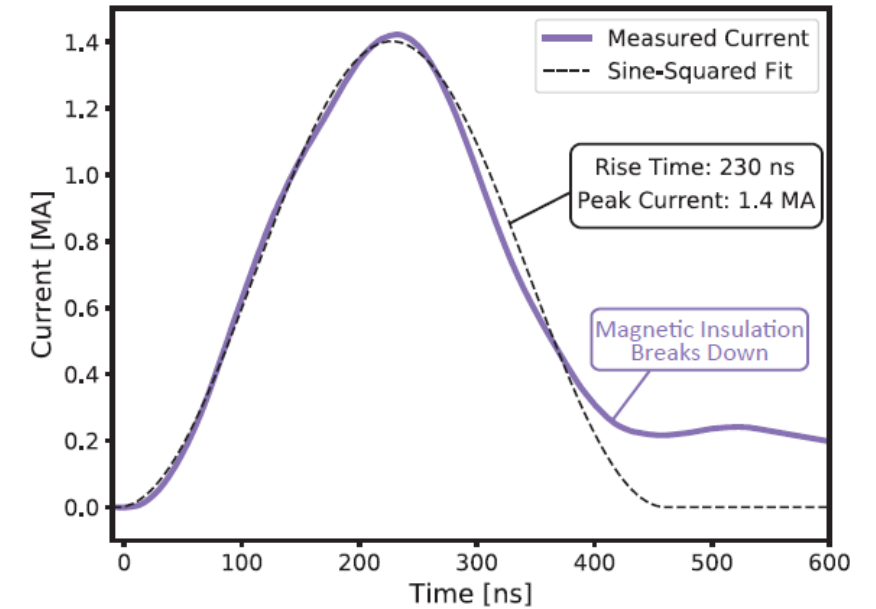
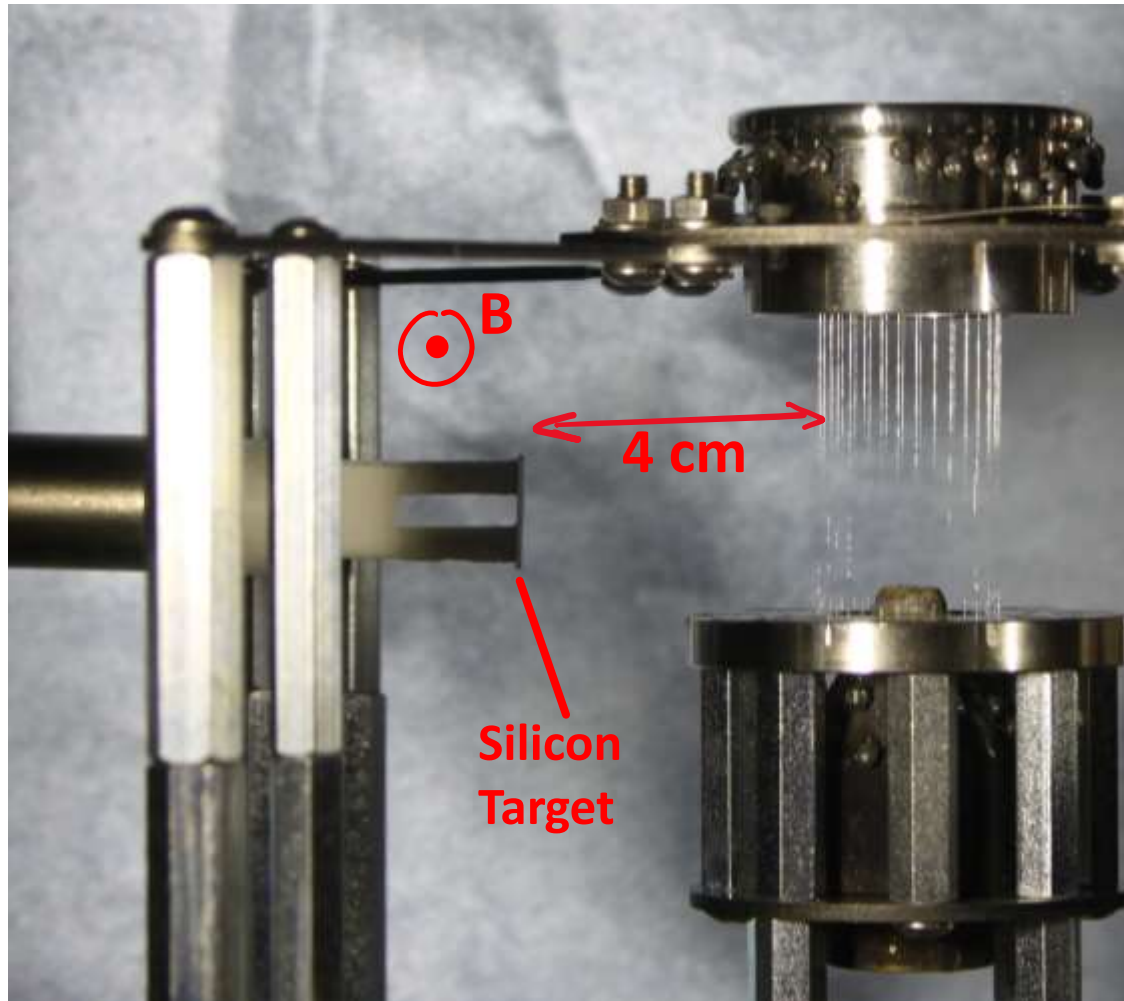
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- Velocity, temperature, & ionisation profiles from Thomson scattering
- Magnetic field profiles from Faraday rotation imaging



1.4 MA, 240 ns Current Pulse

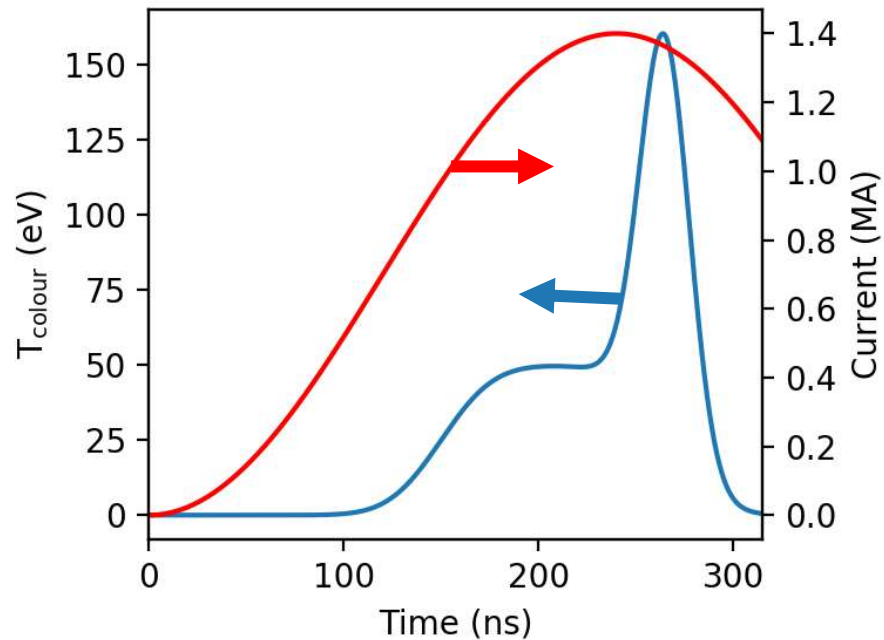
X-Ray Pulse  $\sim$  1 TW



1.4 MA, 240 ns Current Pulse

X-Ray Pulse  $\sim$  1 TW

# Spectral Character of Radiation at $\sim 1$ MA Level

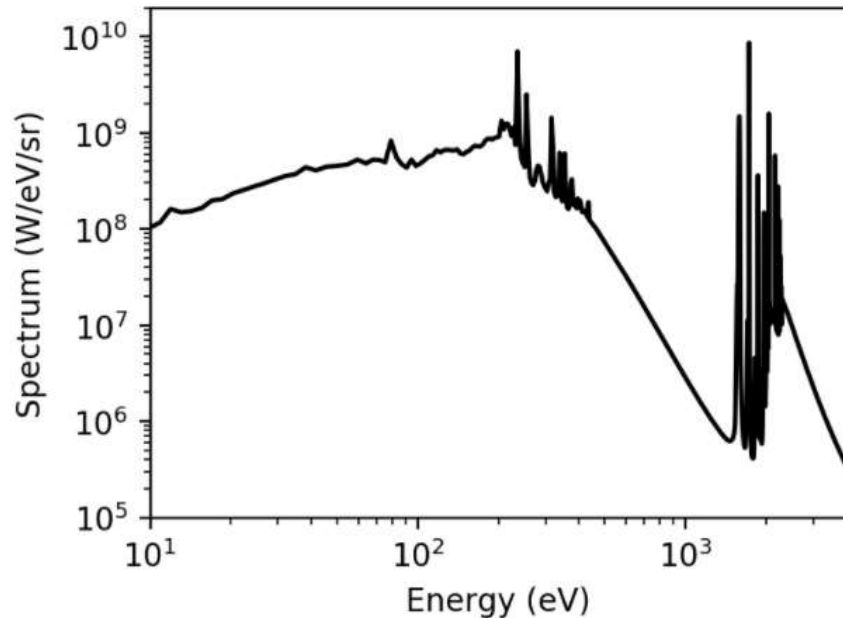


- **Precursor (pre-pulse):**

- Colder spectral character ( $T_c \sim 50$  eV)
- Radiates  $\sim 400$  J in total
- Time duration  $\sim 100$  ns

- **Implosion:**

- Emitted radiation  $\sim 15$  kJ over  $\sim 30$  ns
- Estimate  $T_c \sim 150$  eV



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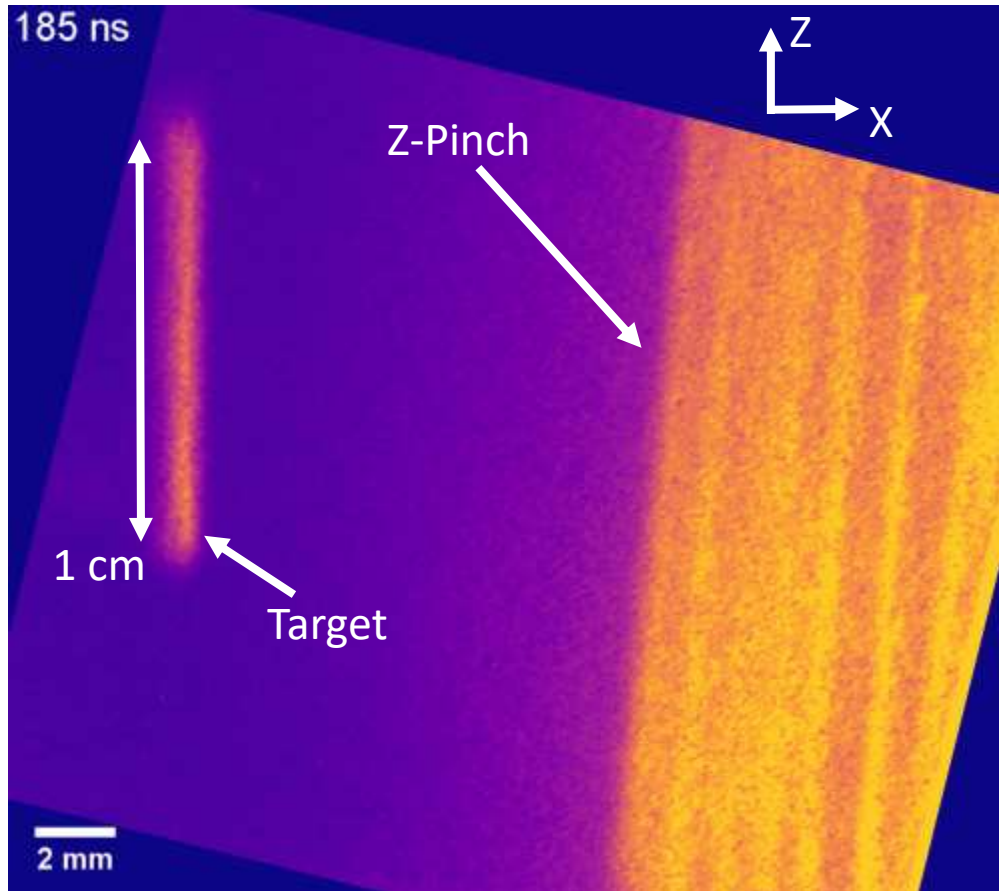
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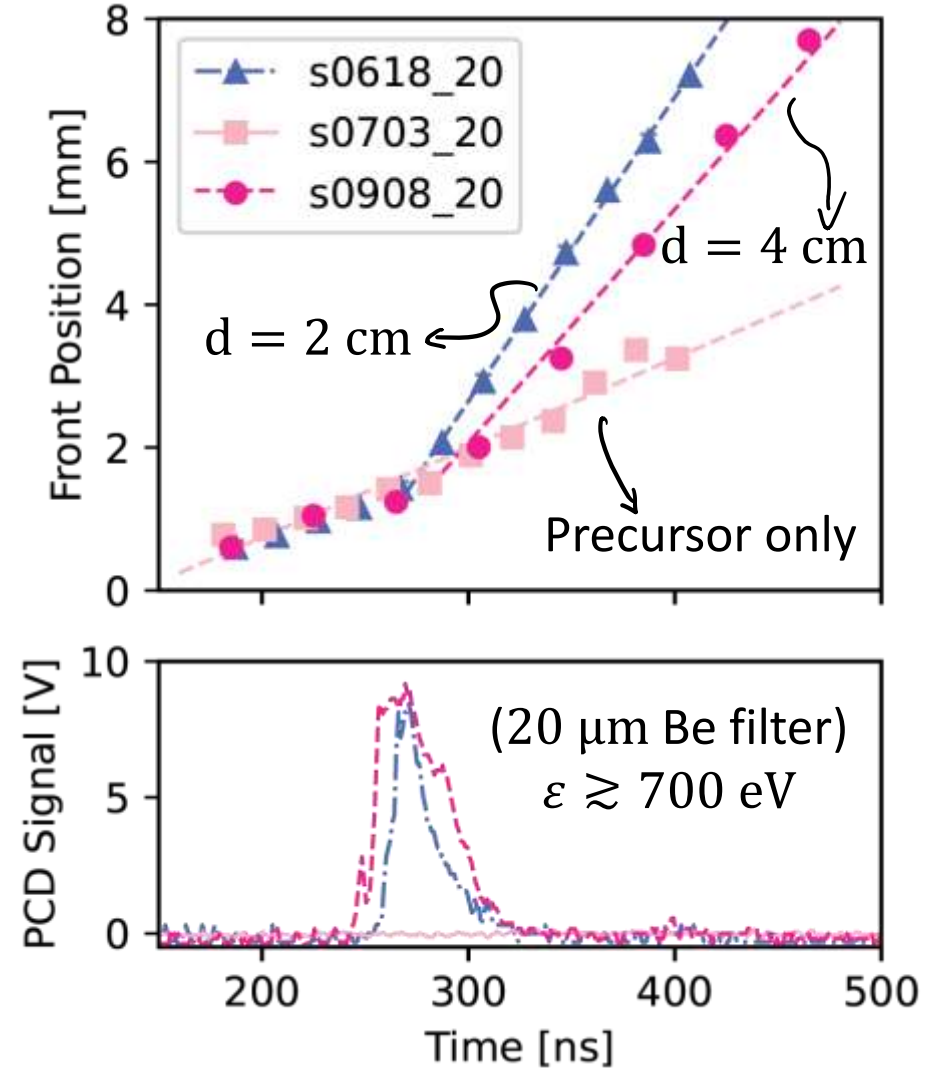


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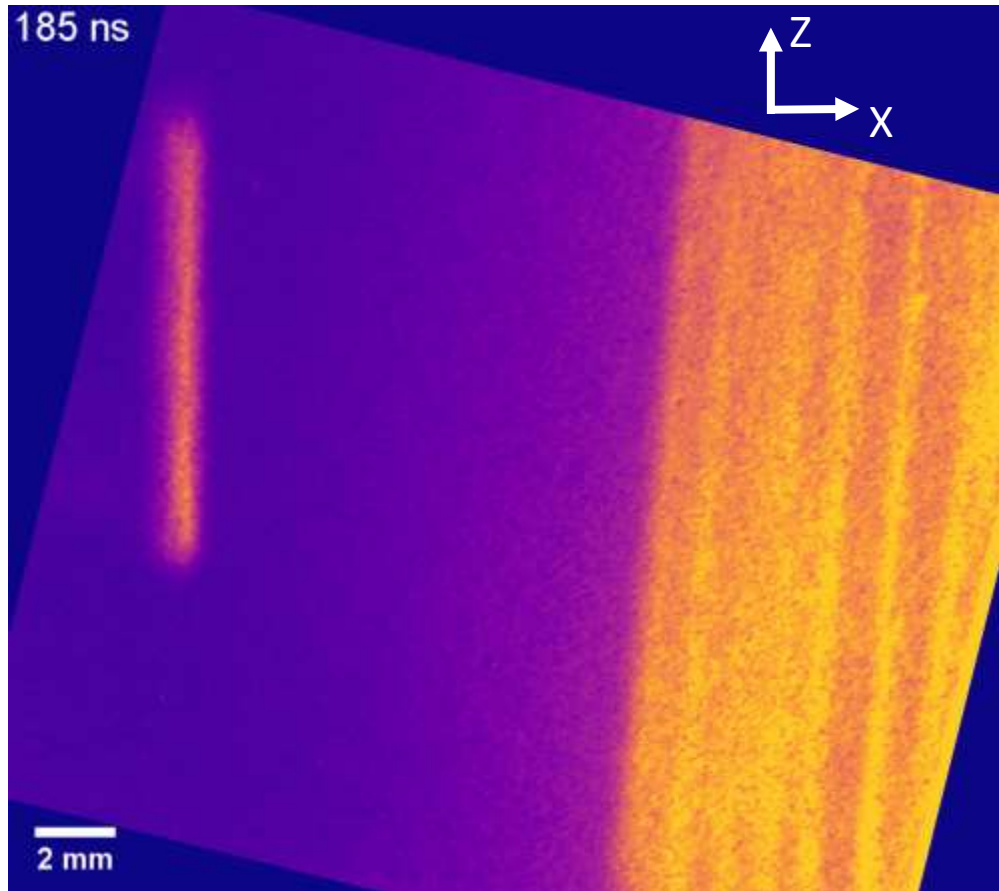
# Optical self emission images [qualitative dynamics]



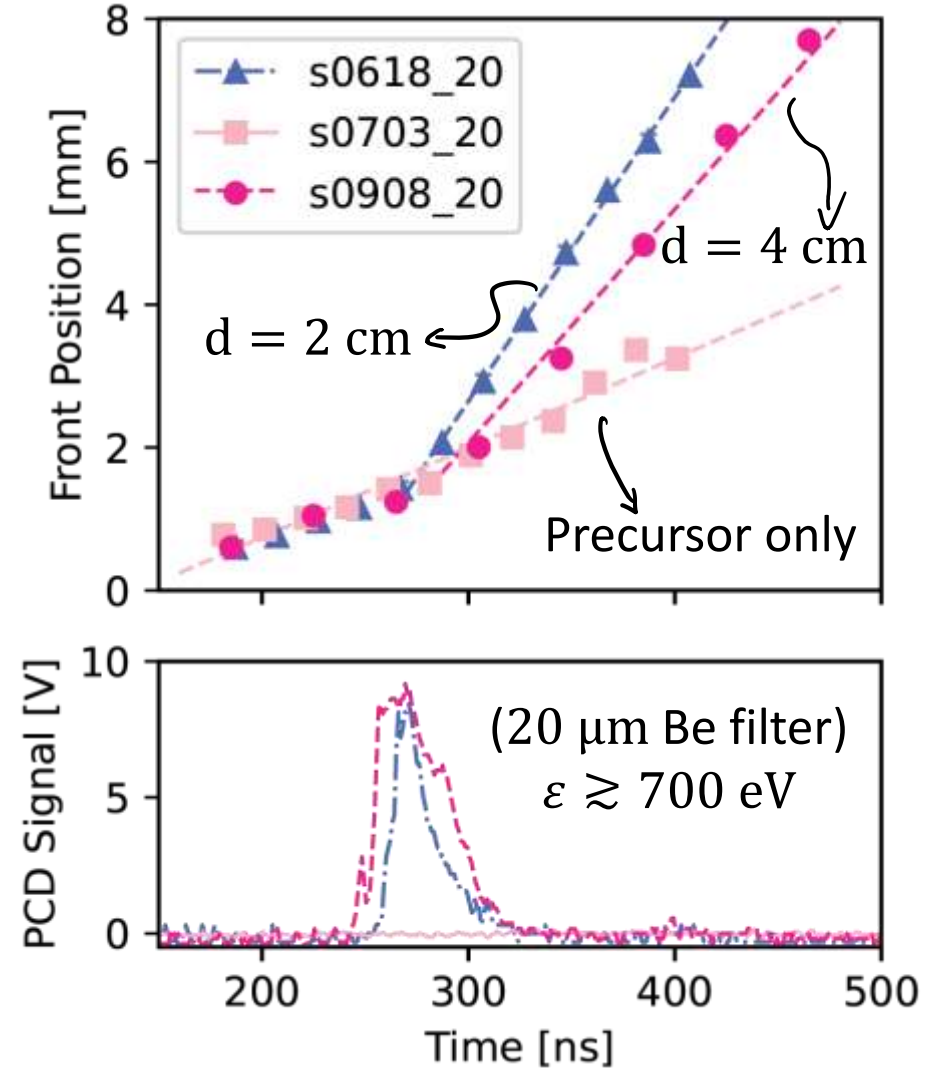
Self emission images [ $600 \lesssim \lambda \lesssim 900 \text{ nm}$ ]



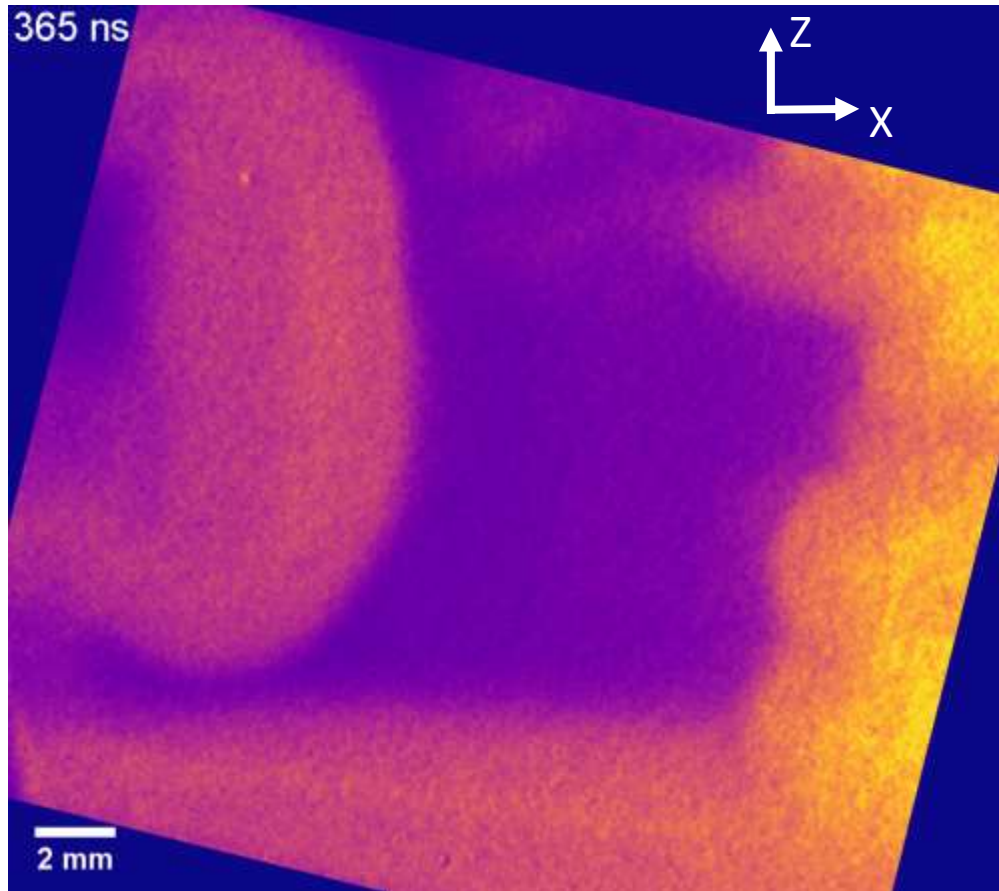
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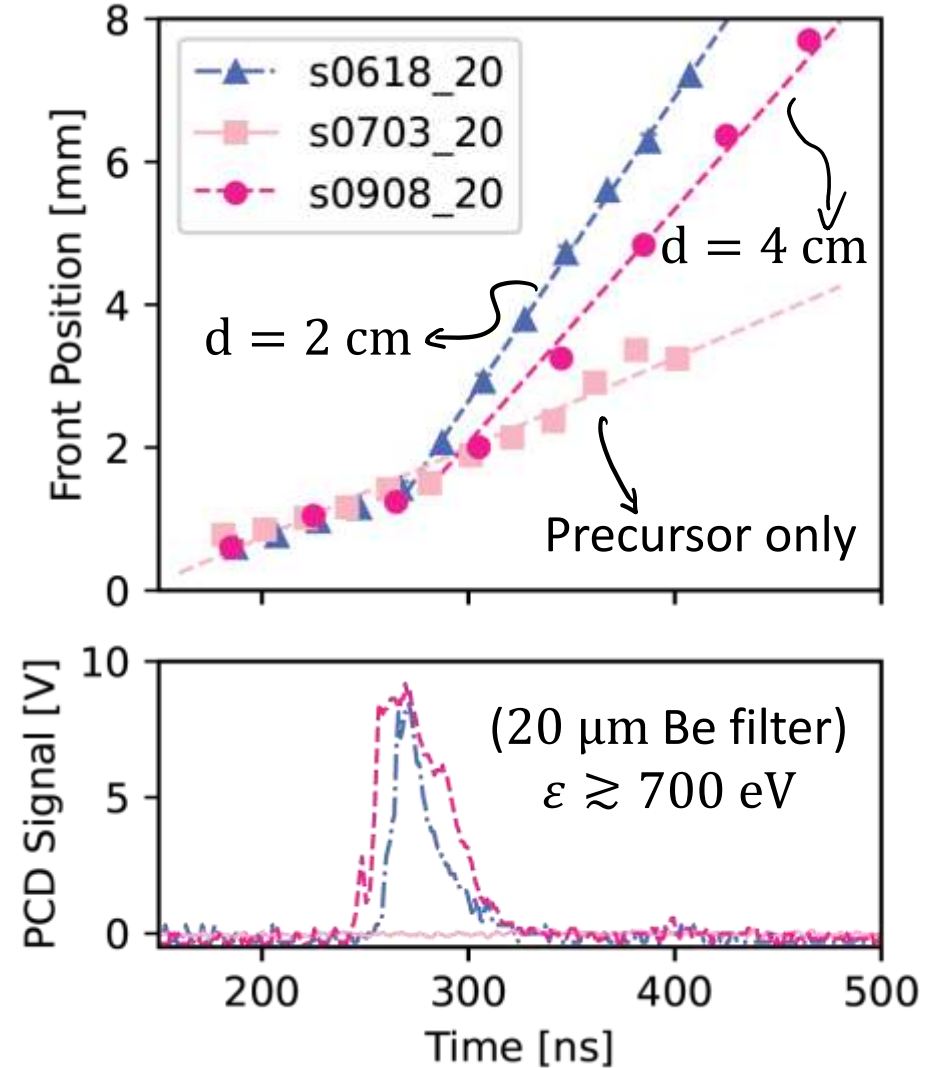
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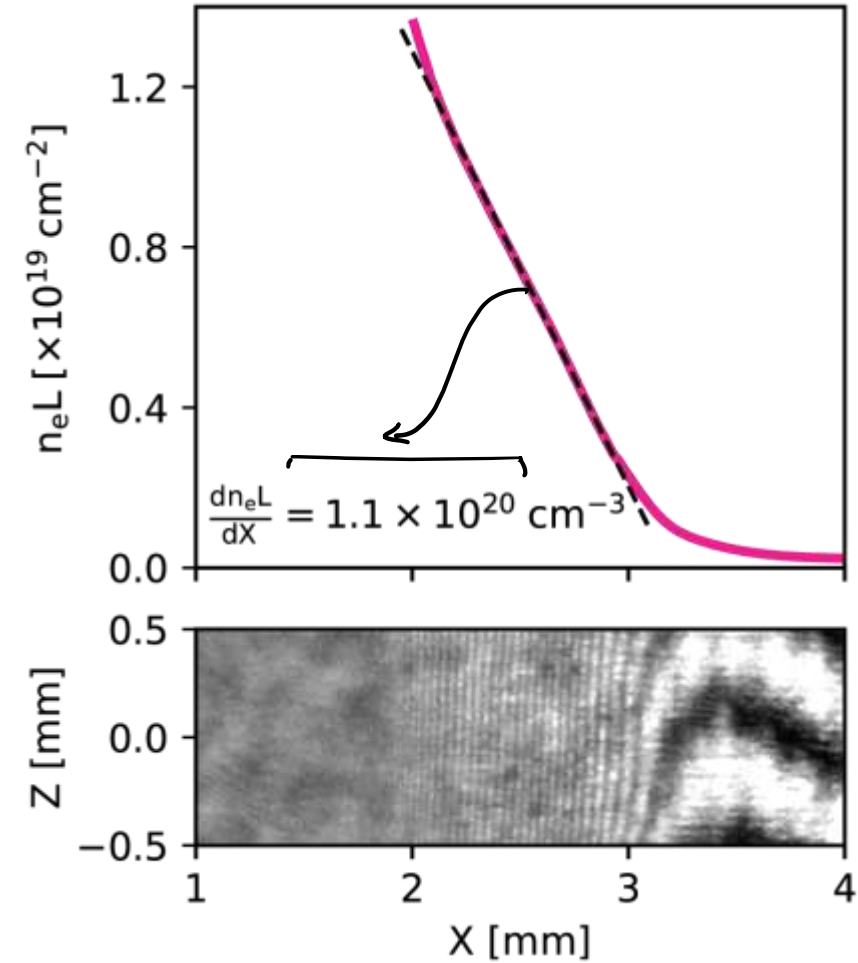
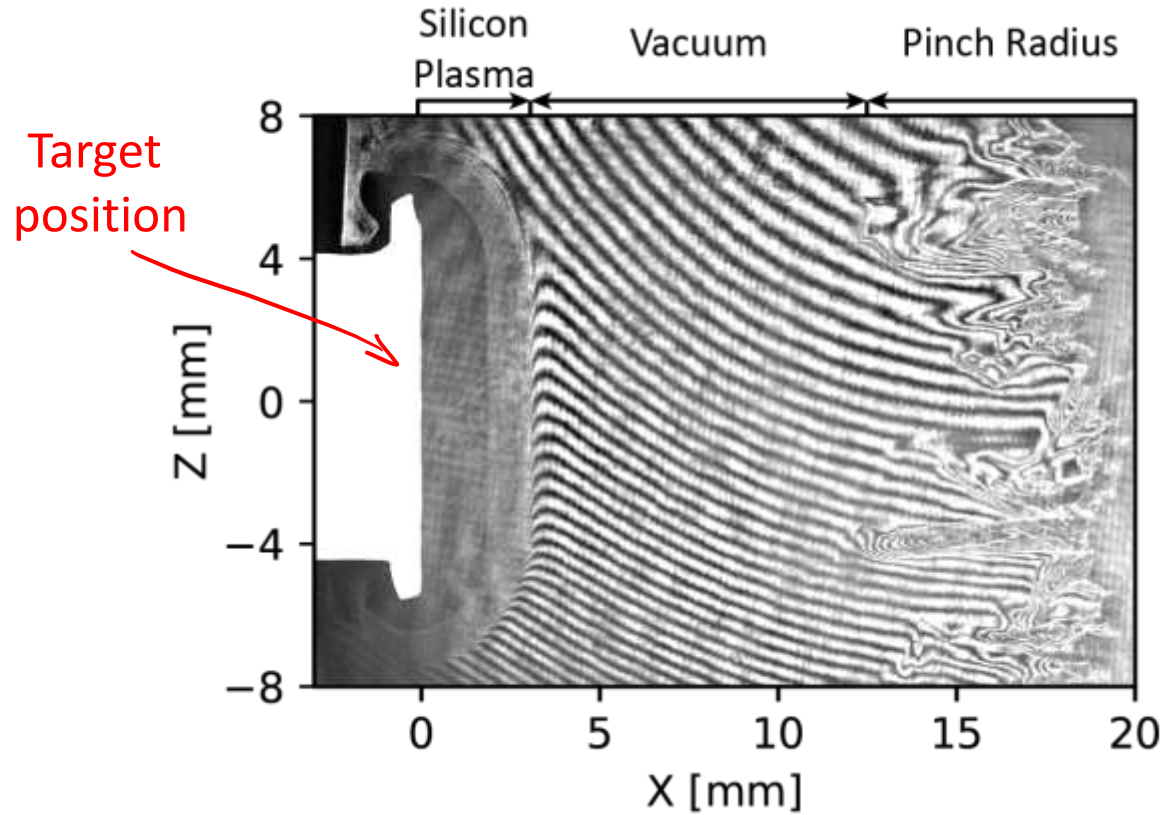
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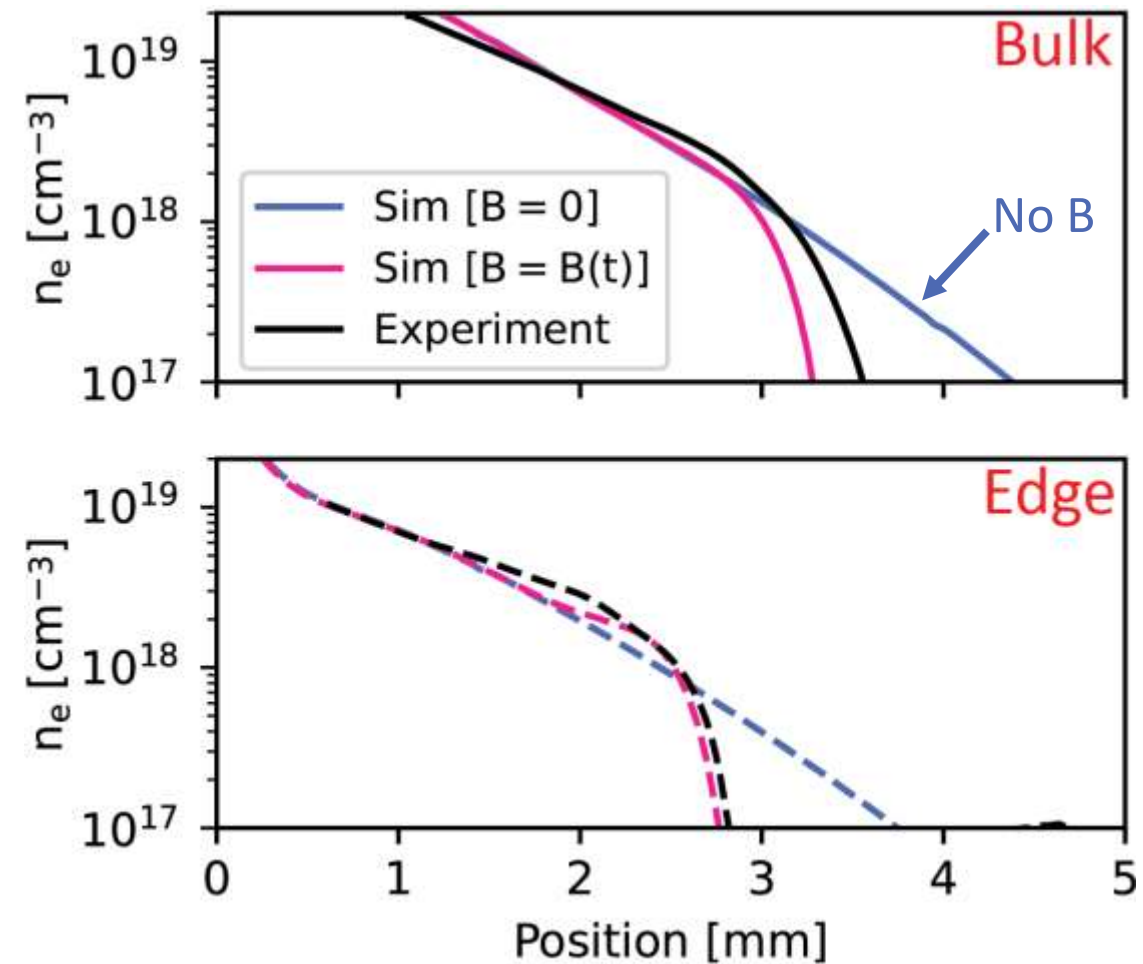
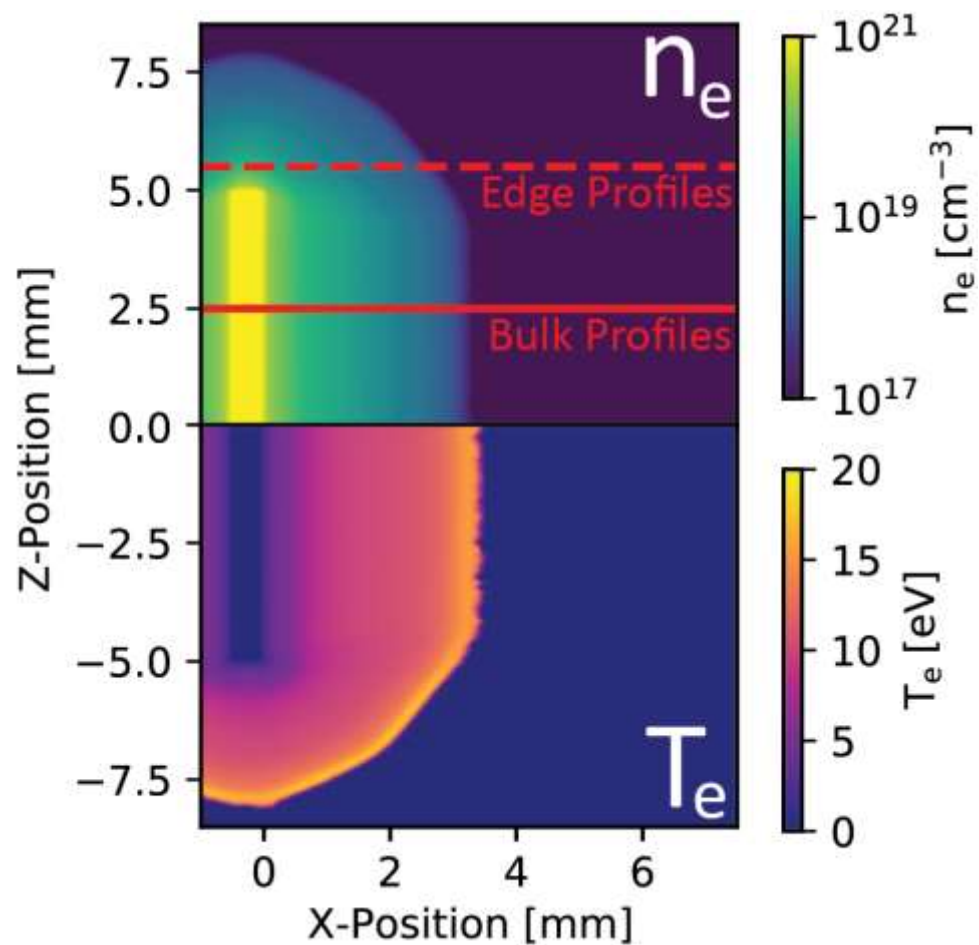


# Interferometry [line integrated electron density]



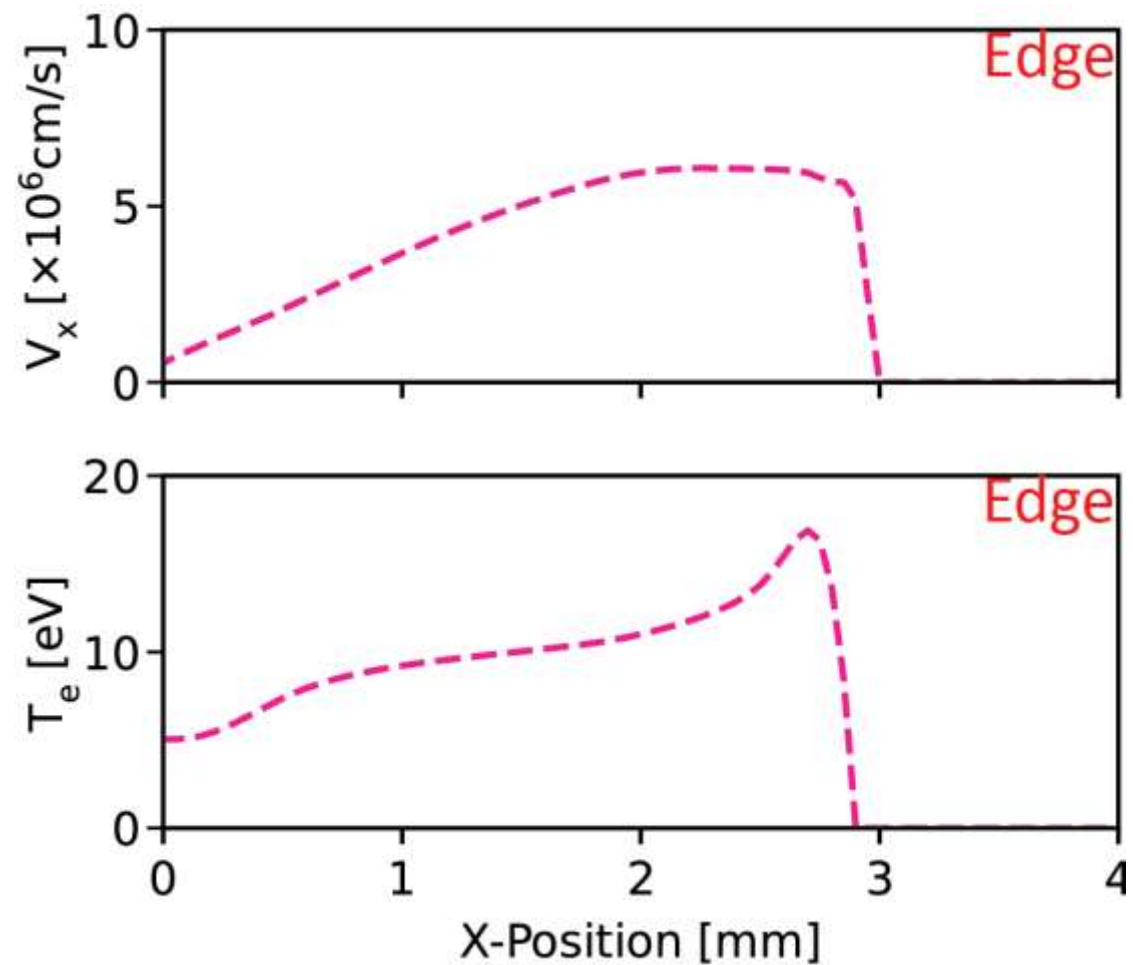
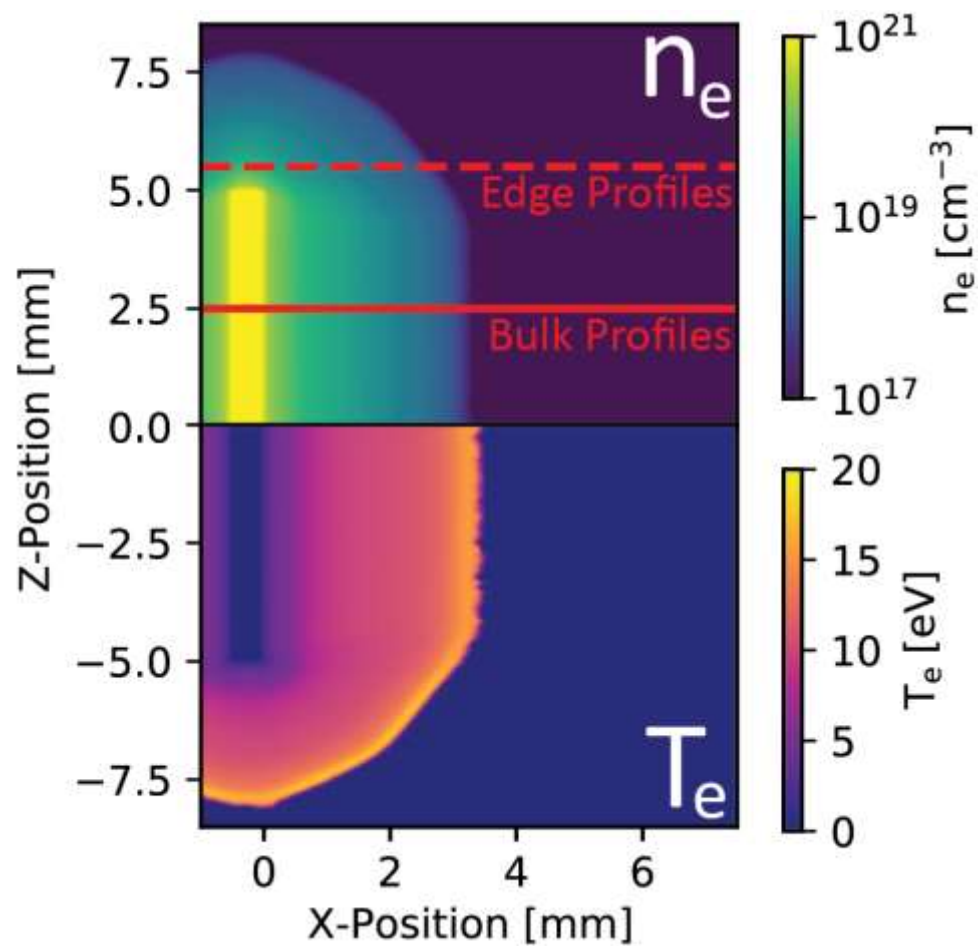
- Interferogram captured at  $t = 320 \text{ ns}$
- Smooth  $\sim 1\text{D}$  expansion profile confirmed by orthogonal laser probing

## Radiative MHD simulations [Chimera]



Density profile is affected by B-Field.

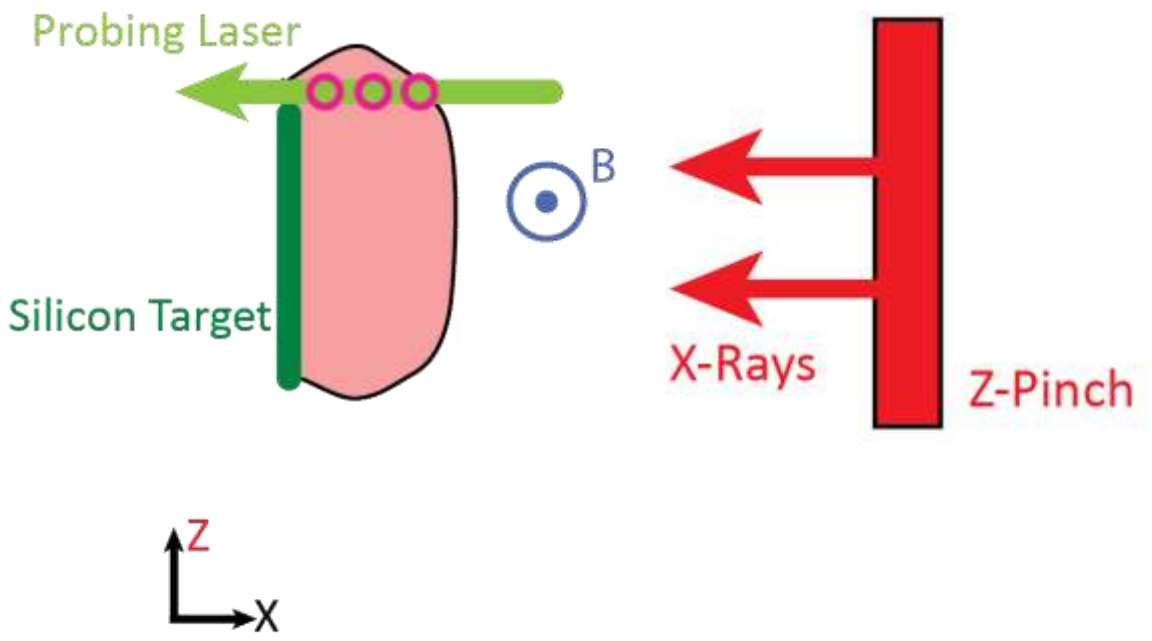
## Radiative MHD simulations [Chimera]



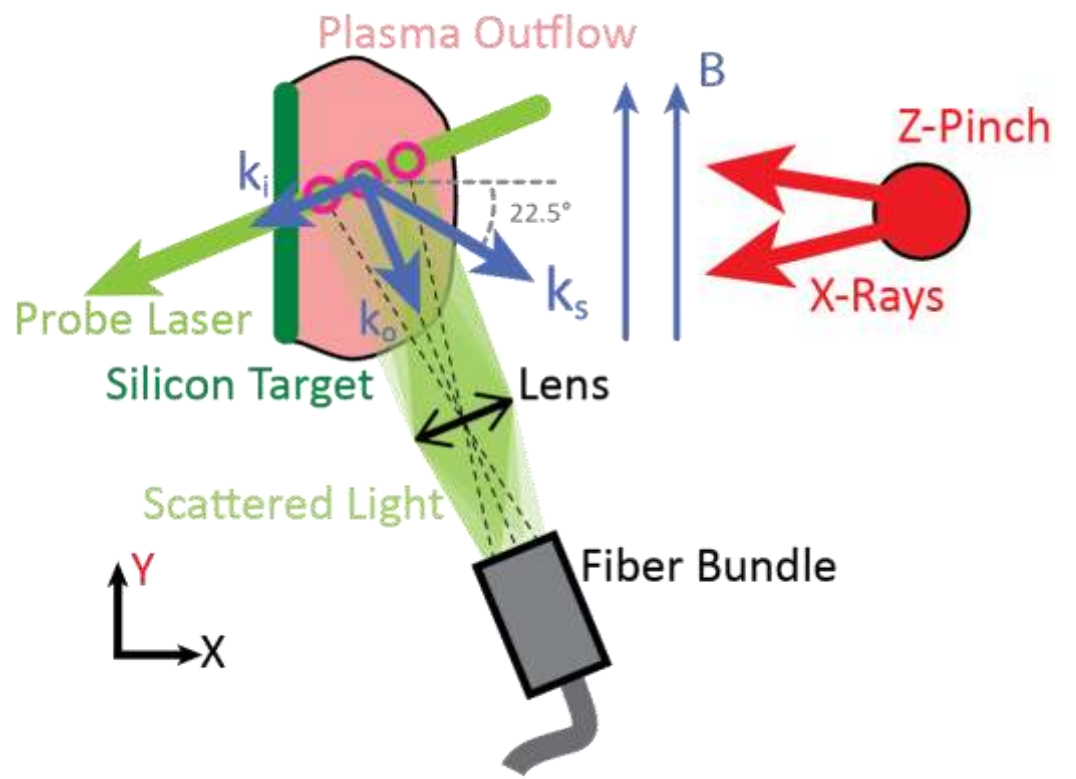
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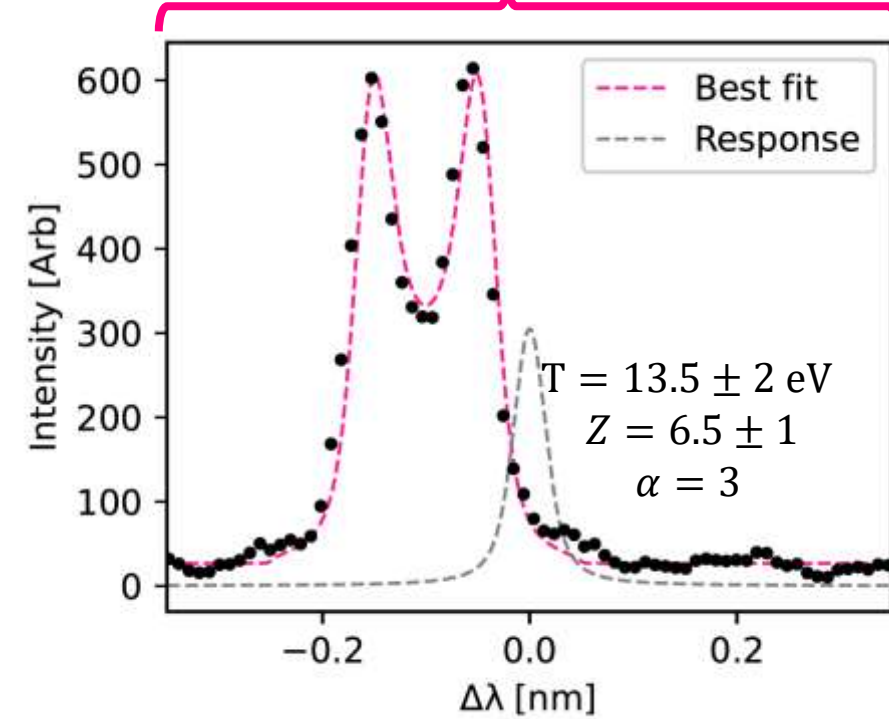
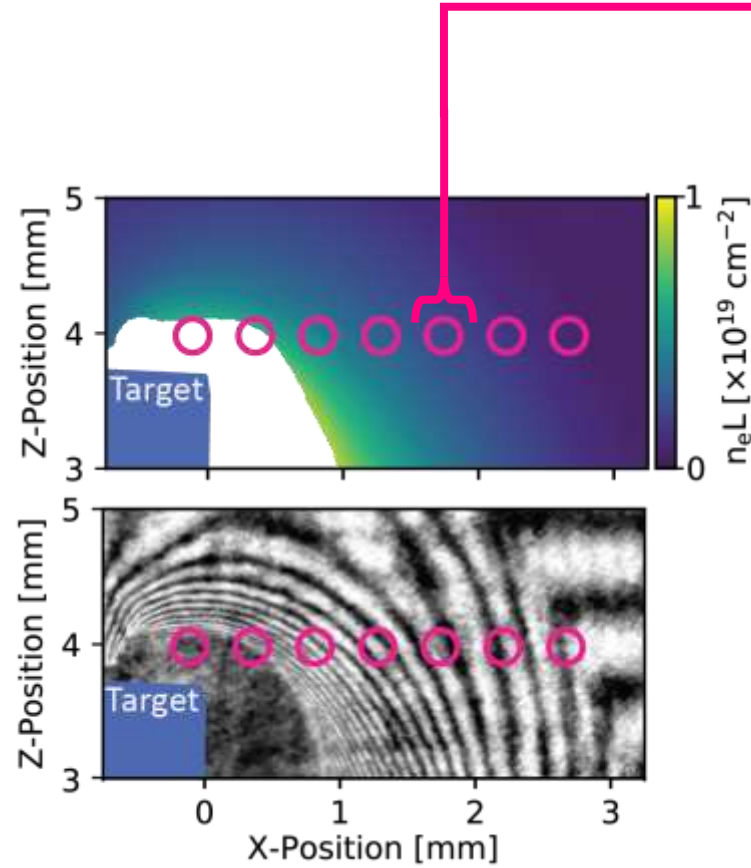
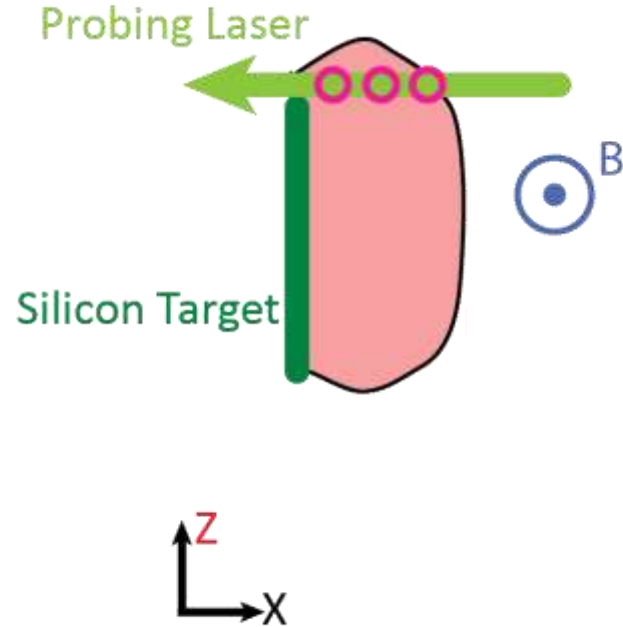
# Thomson scattering [localised diagnosis of $T, V, Z$ ]



Side-On (X-Z plane) view of the experiment



End-On (X-Y plane) view of the experiment

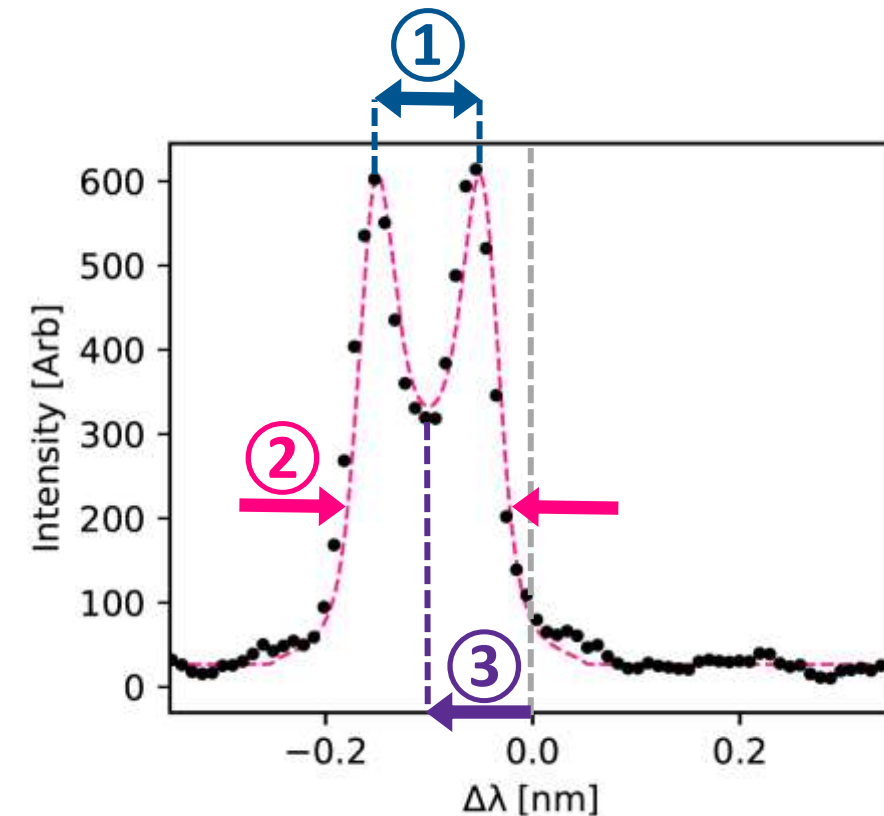


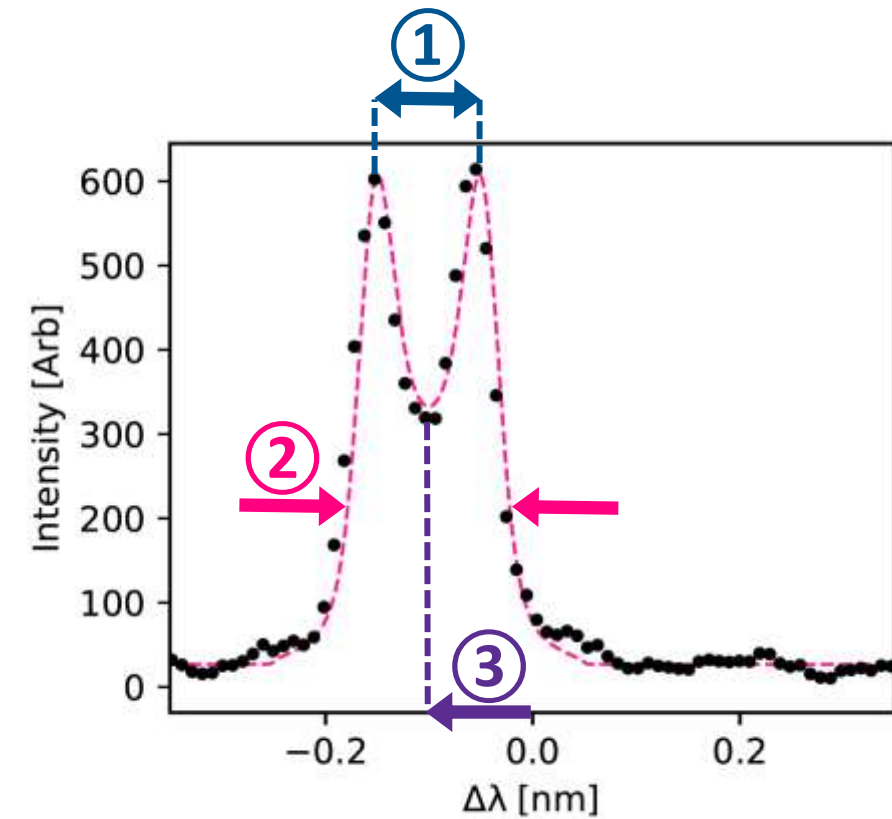
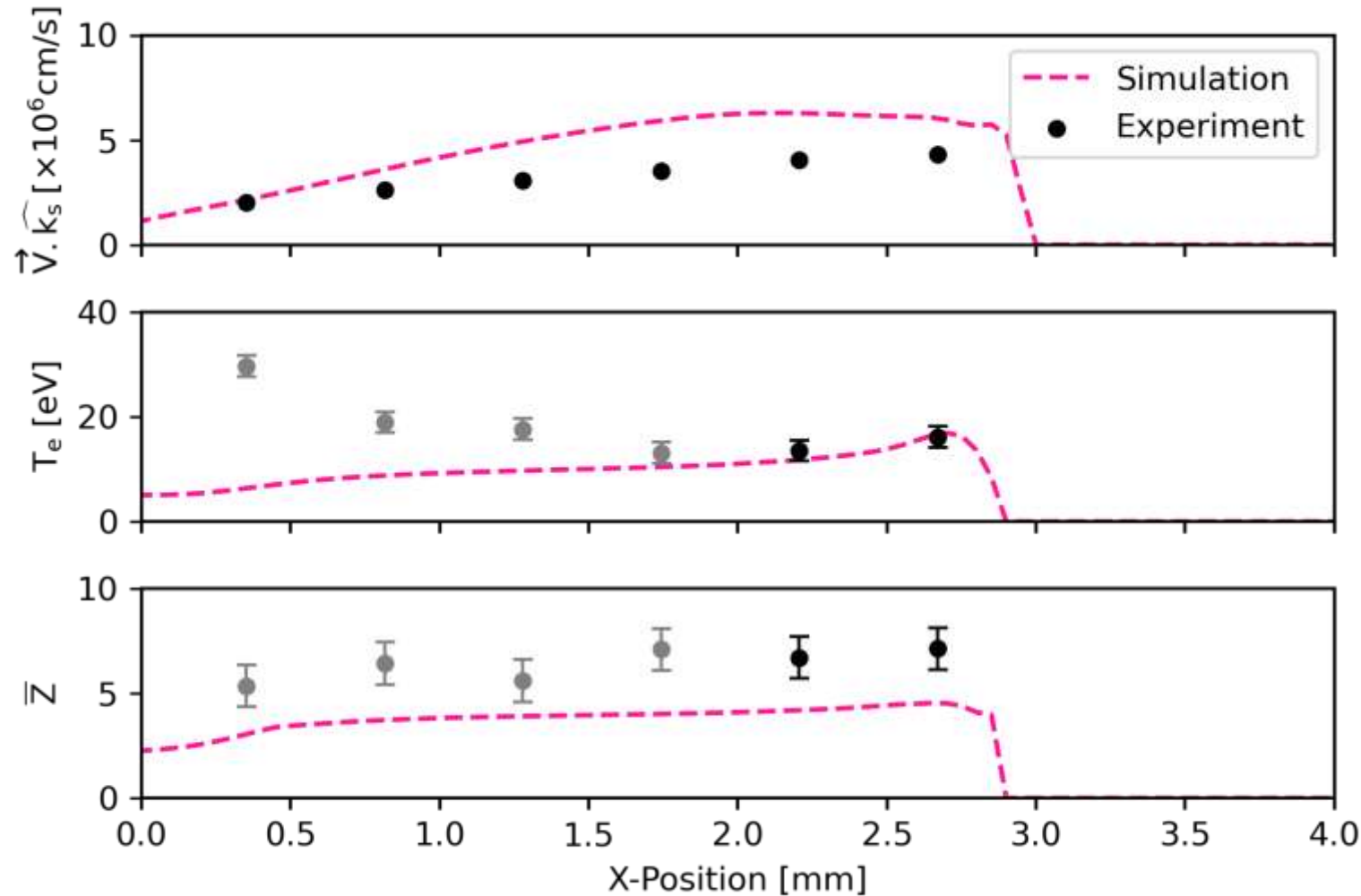
- ①: Ion Acoustic peak **separation** depends on  $\bar{Z} \times T_e$
- ②: Feature **width** depends on  $n_e$ ,  $T_i$ , and spectral response
- ③: Doppler **shift** from probe wavelength depends on  $\vec{V} \cdot \hat{k}_s$

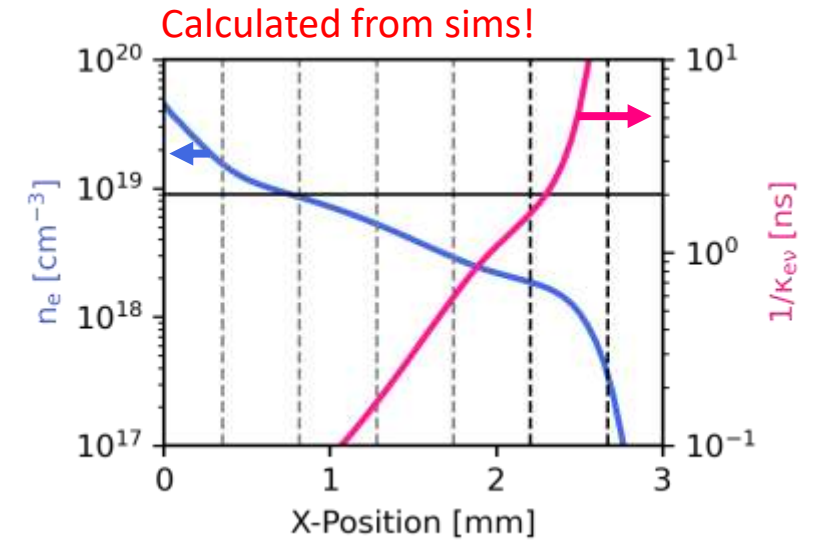
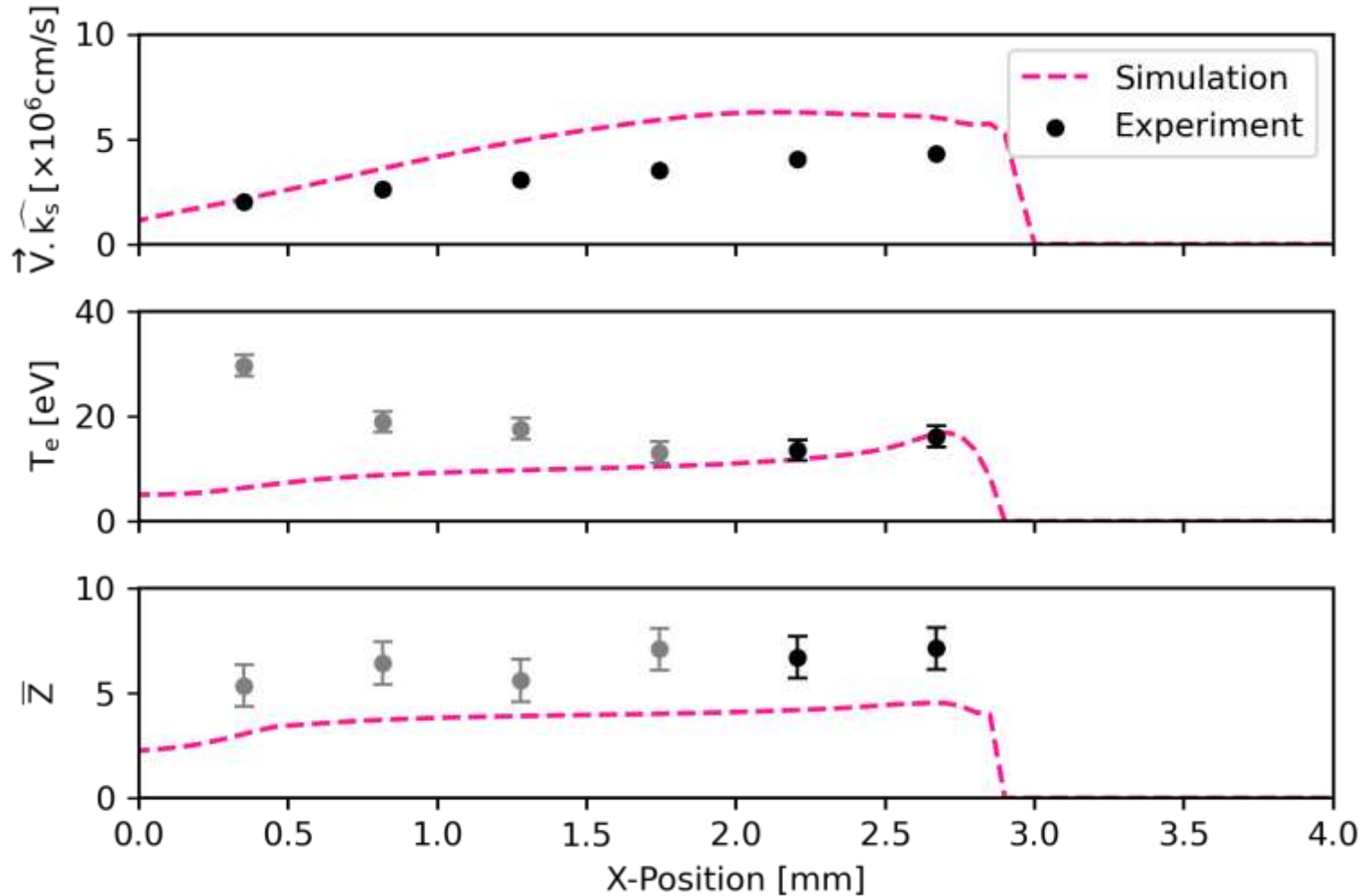
Enforced  $T_e = T_i$ , and allowed  $\bar{Z}$  to vary ( $\tau_{ei} \lesssim 1$  ns).

Convolved calculated spectra with measured spectral response.

Constrained value of  $n_e$  from (near simultaneous) interferometry.



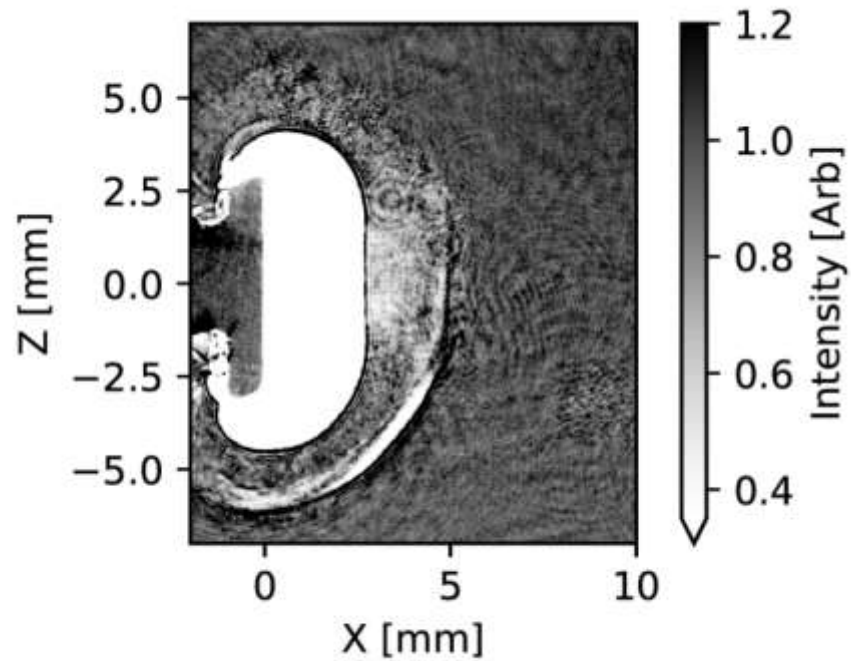




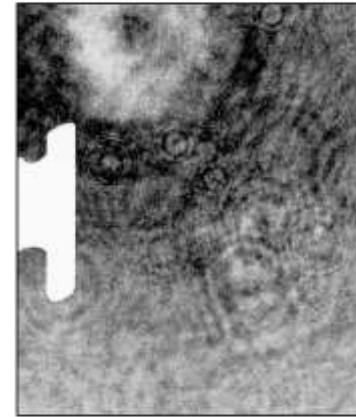
$$\kappa_{ve} \propto \frac{Zn_e^2 \ln(\Lambda) T_e^{-\frac{3}{2}}}{\sqrt{\omega^2 \left(1 - \frac{\omega_p^2}{\omega^2}\right)}}$$

N. R. L. plasma physics formulary

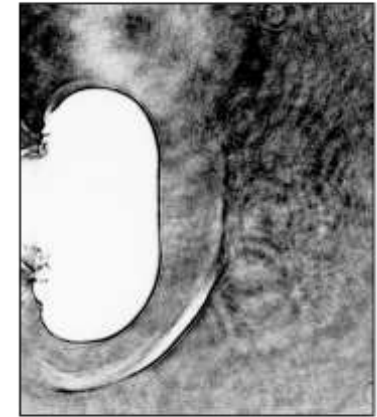
# Measurement of Inverse Bremsstrahlung Coefficient



$$I = I_0 e^{-\kappa_{ve} x} \Rightarrow \kappa_{ve} = \frac{-\ln(I/I_0)}{x}$$



Background Image ( $I_0$ )

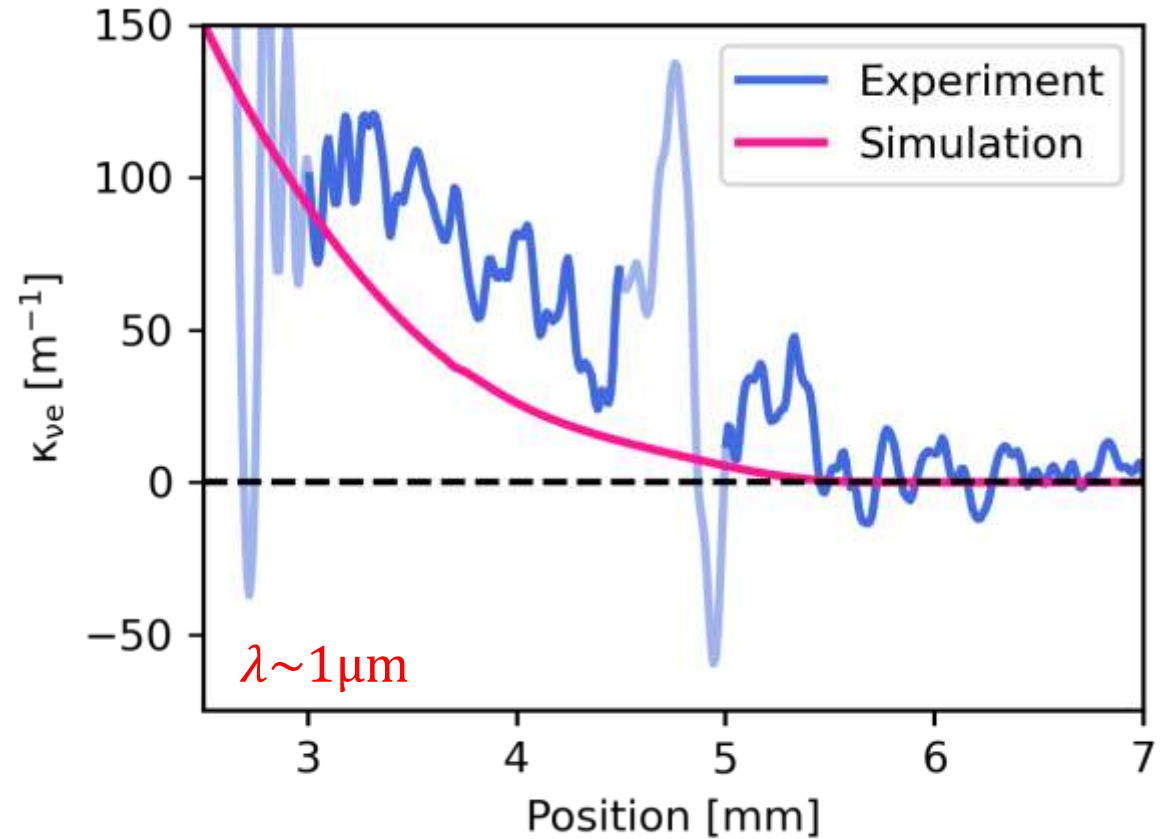
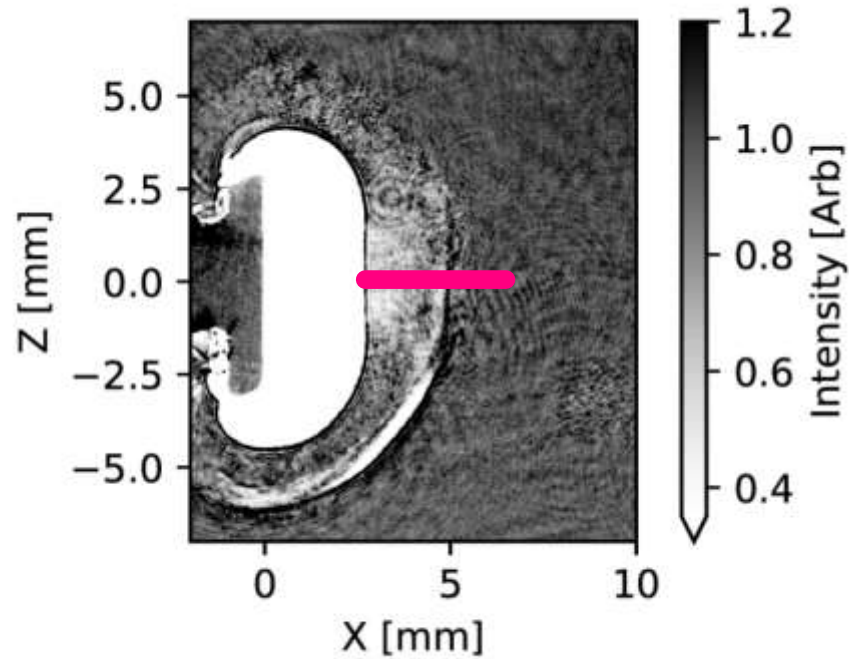


Shot Image ( $I$ )

$$\kappa_{ve} \propto \frac{Z n_e^2 \ln(\Lambda) T_e^{-\frac{3}{2}}}{\sqrt{\omega^2 \left(1 - \frac{\omega_p^2}{\omega^2}\right)}}$$

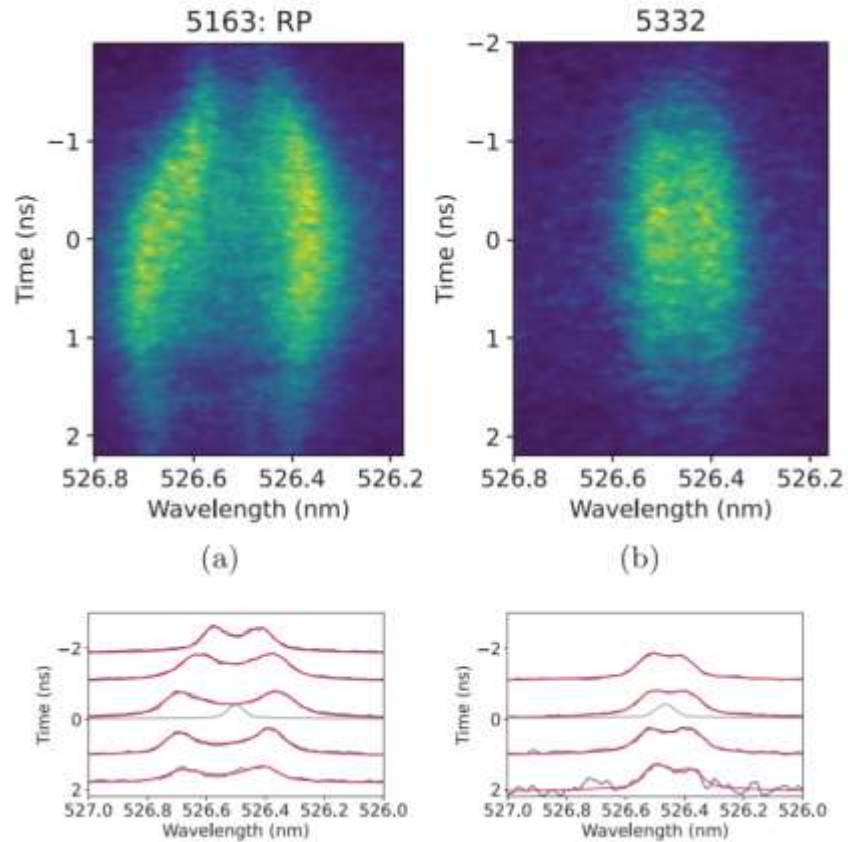
N. R. L. plasma physics formulary

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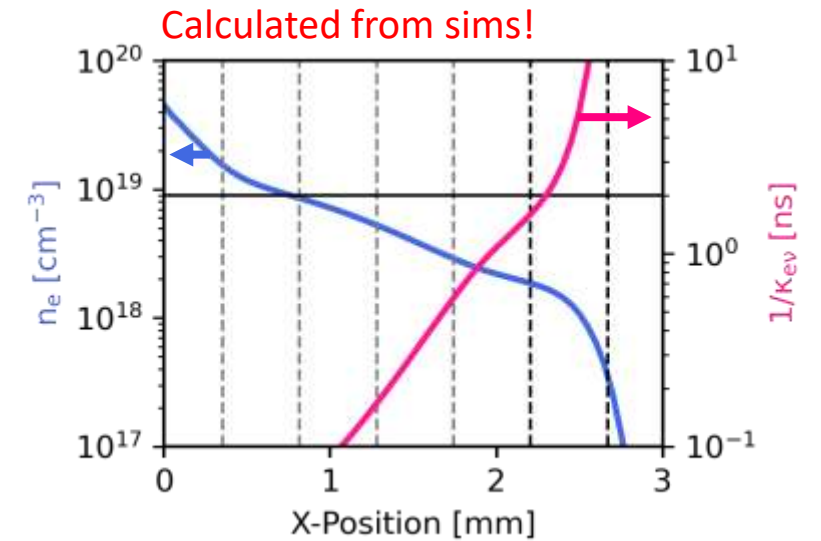


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N. R. L. plasma physics formulary



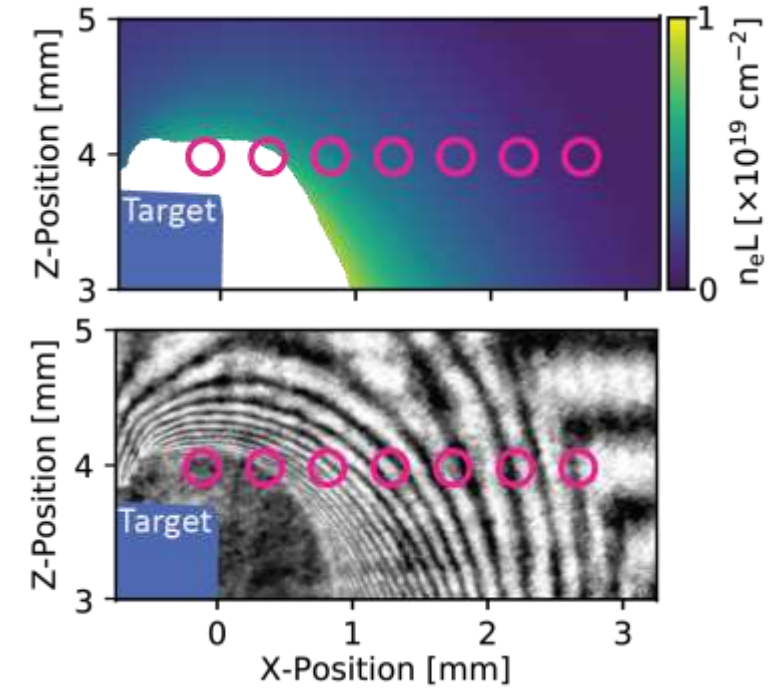
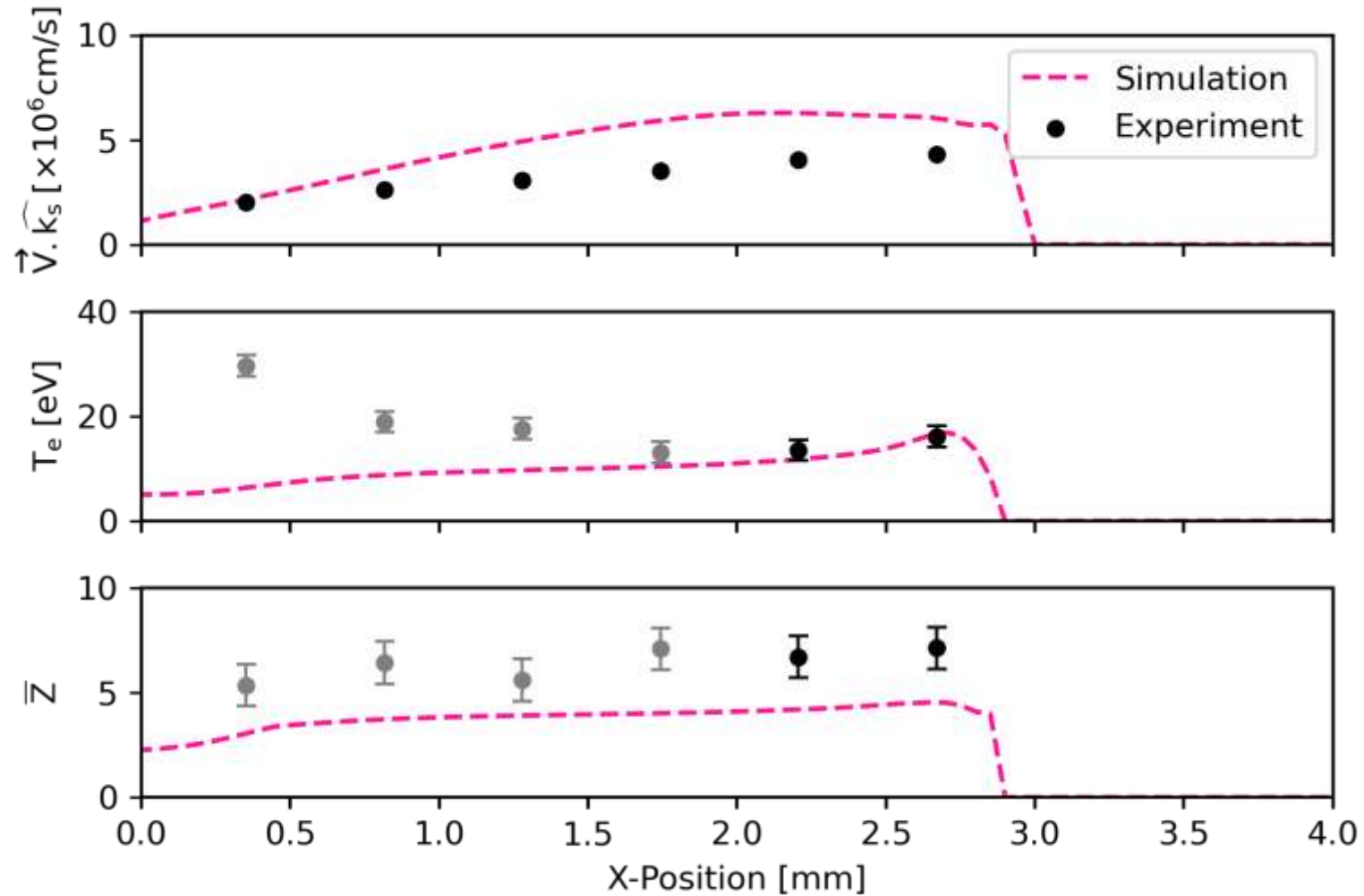
J. T. Banasek *et al* "Current polarity effects on laboratory plasma jets" *Physics of Plasmas* 2021.



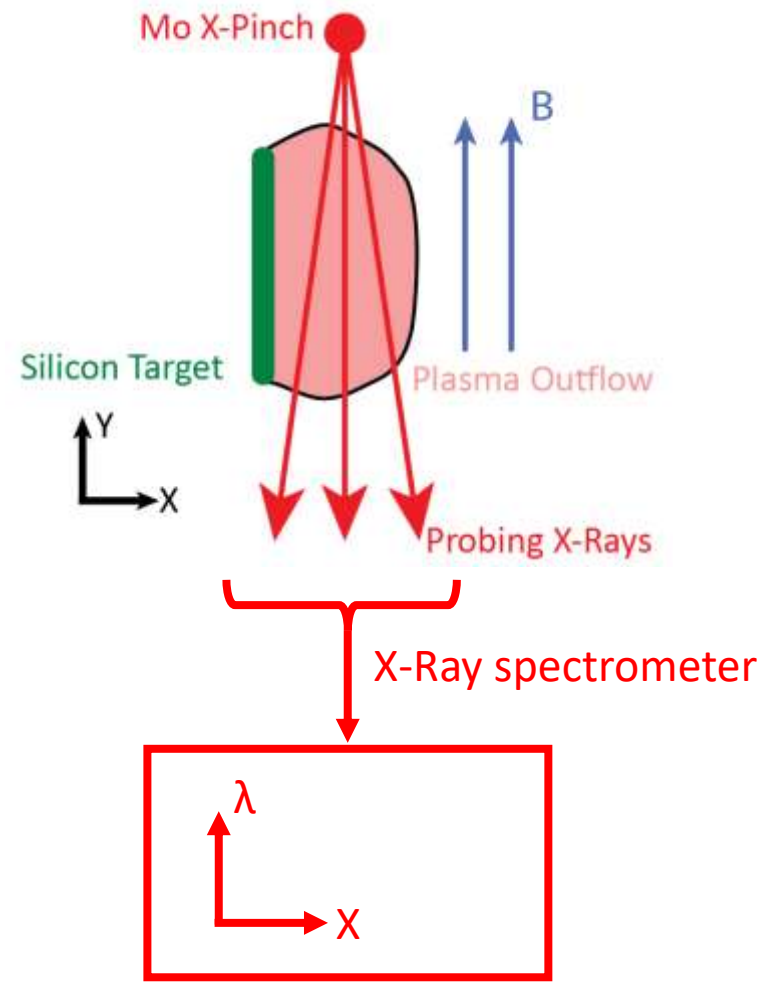
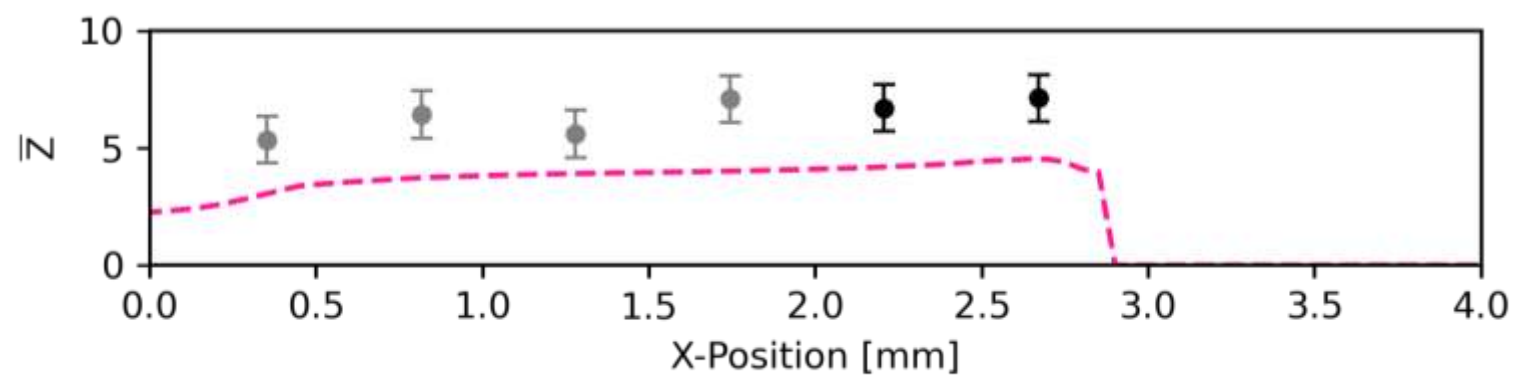
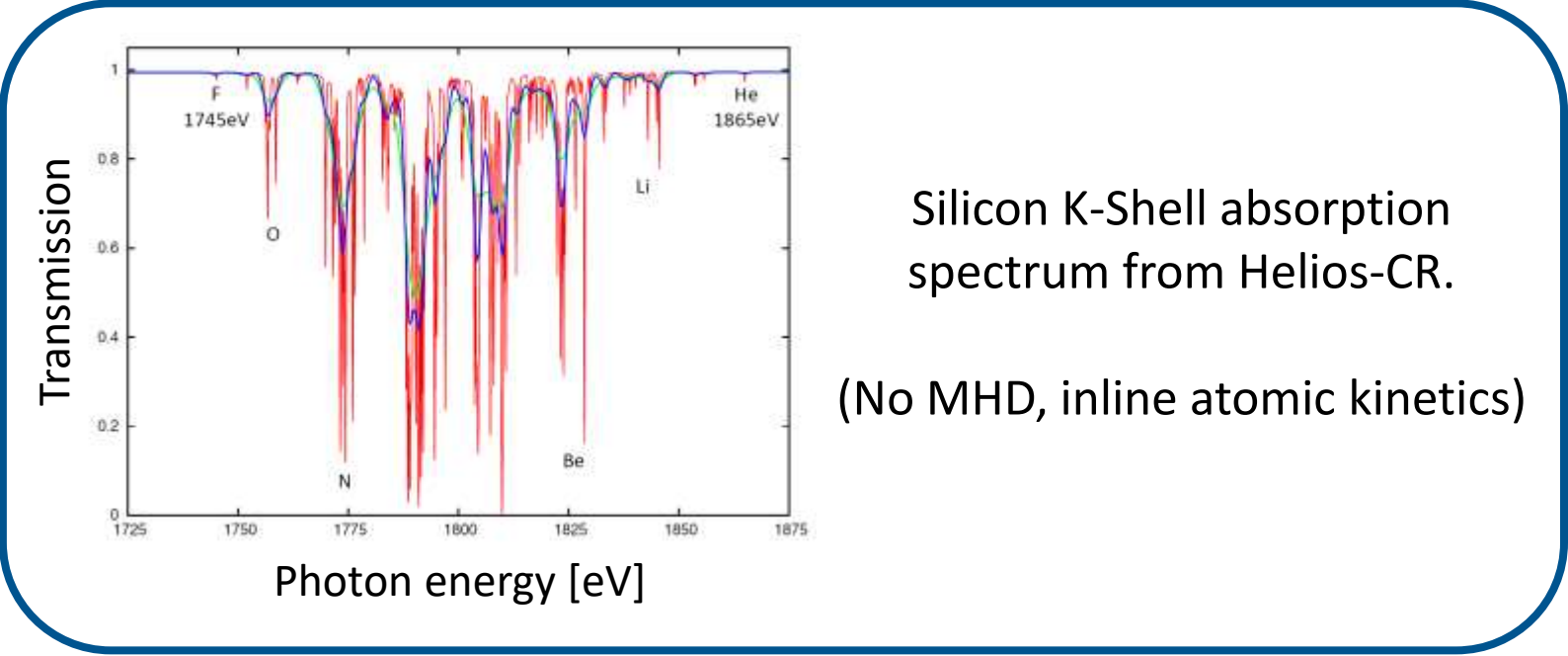
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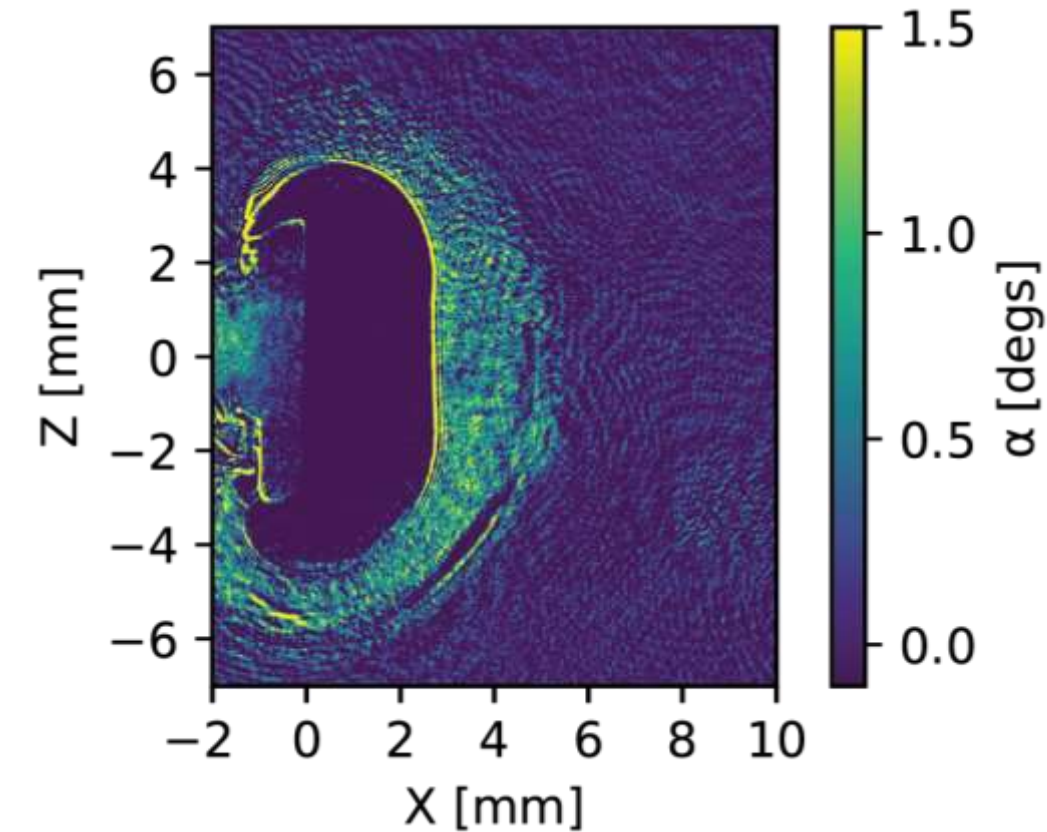




# Future work – Diagnosis of Charge State Distribution



- Discuss X-Ray driver (MAGPIE generator, wire array Z-pinches)
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- Measure rotation applied to laser polarisation:

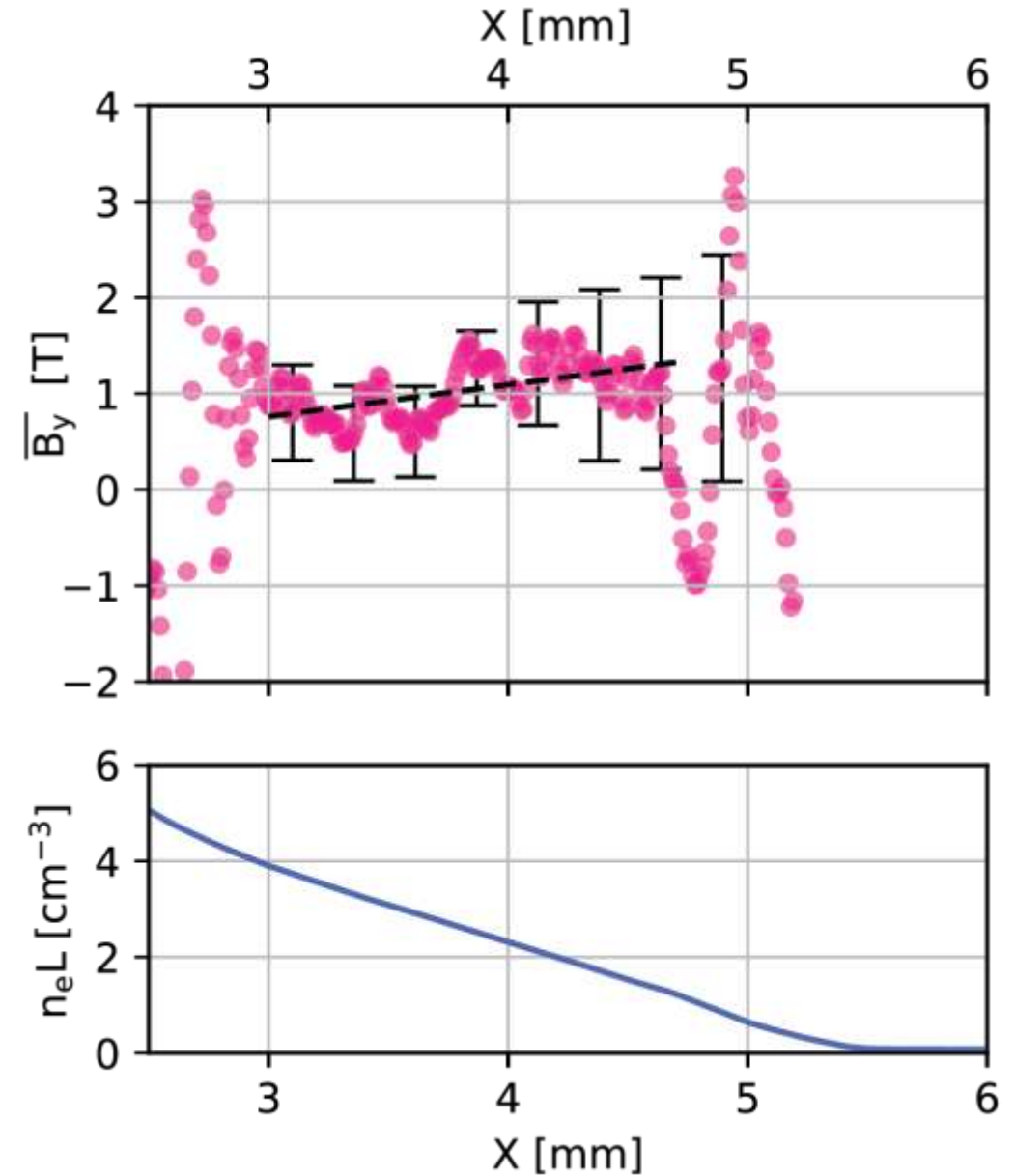
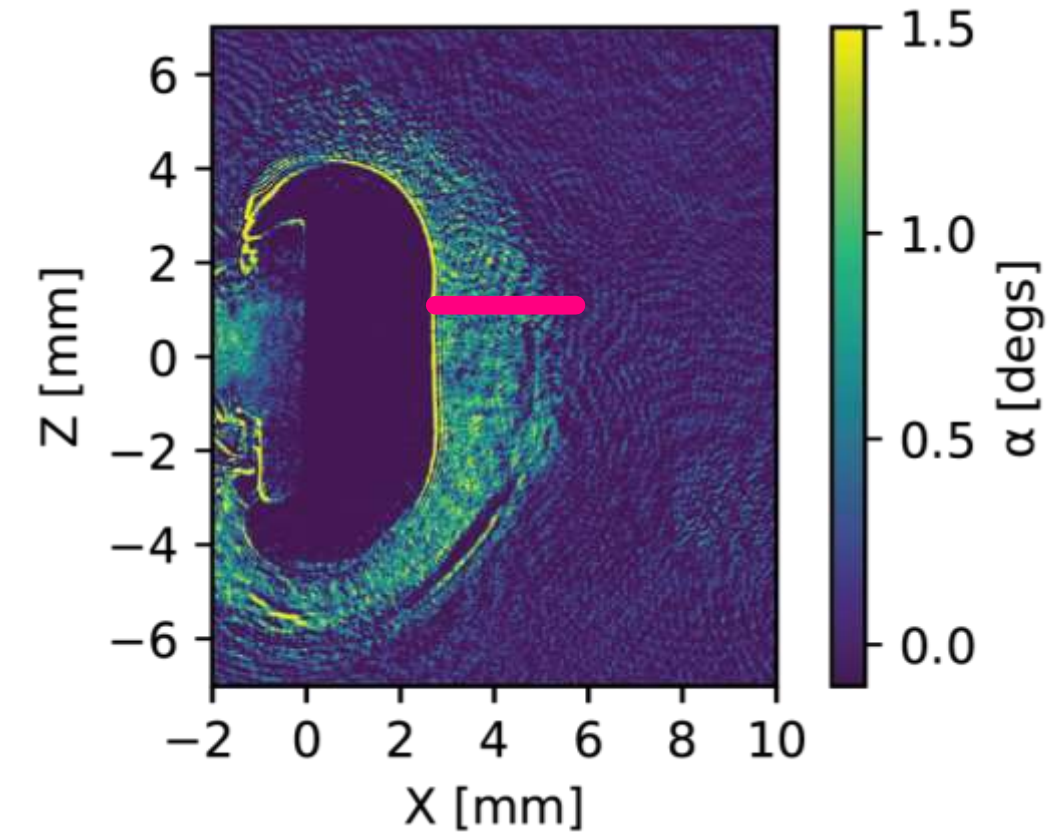
$$\alpha \propto \lambda^2 \int n_e \vec{B} \cdot d\vec{y}$$

- Obtain interferometry along same line of sight:

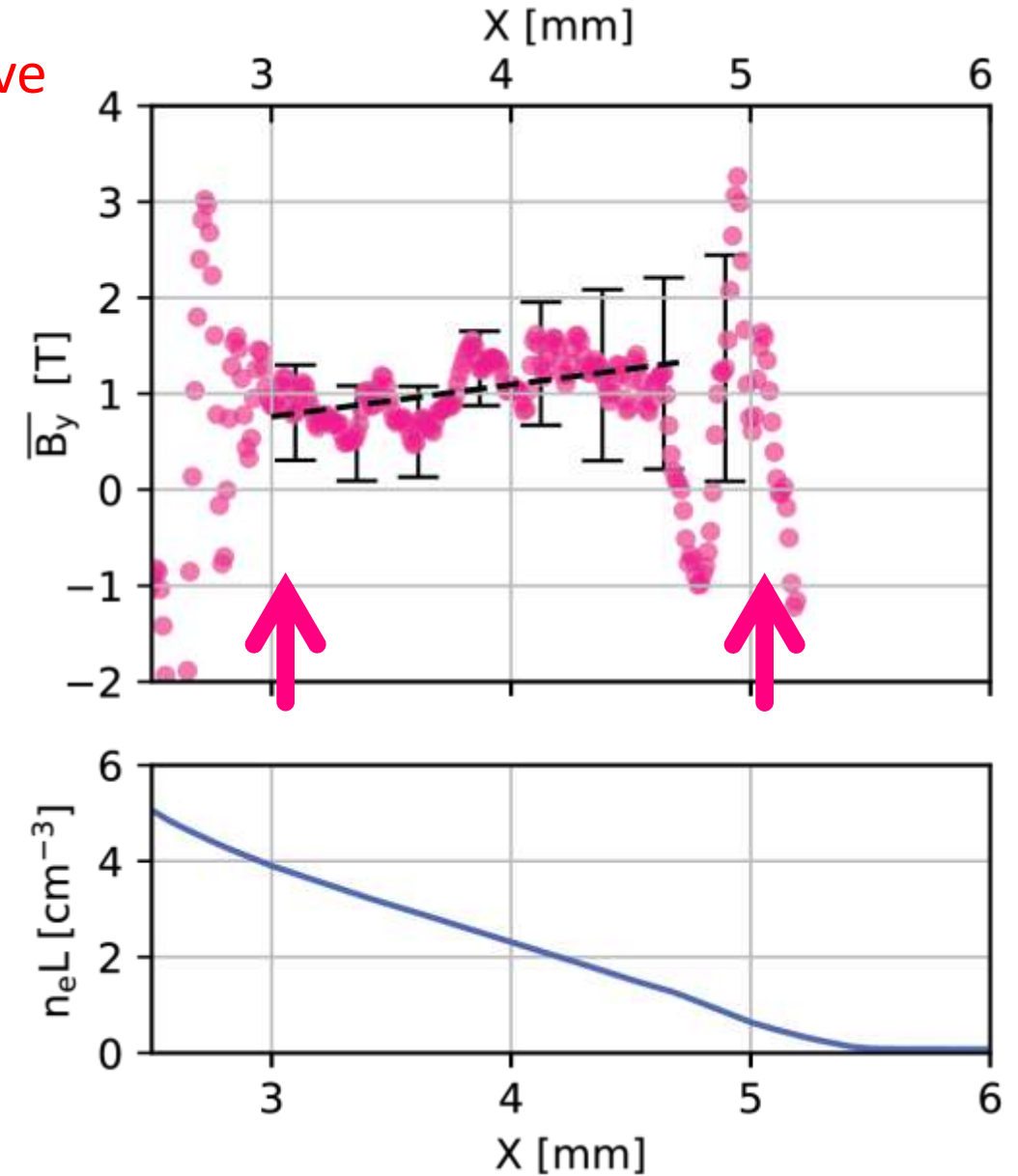
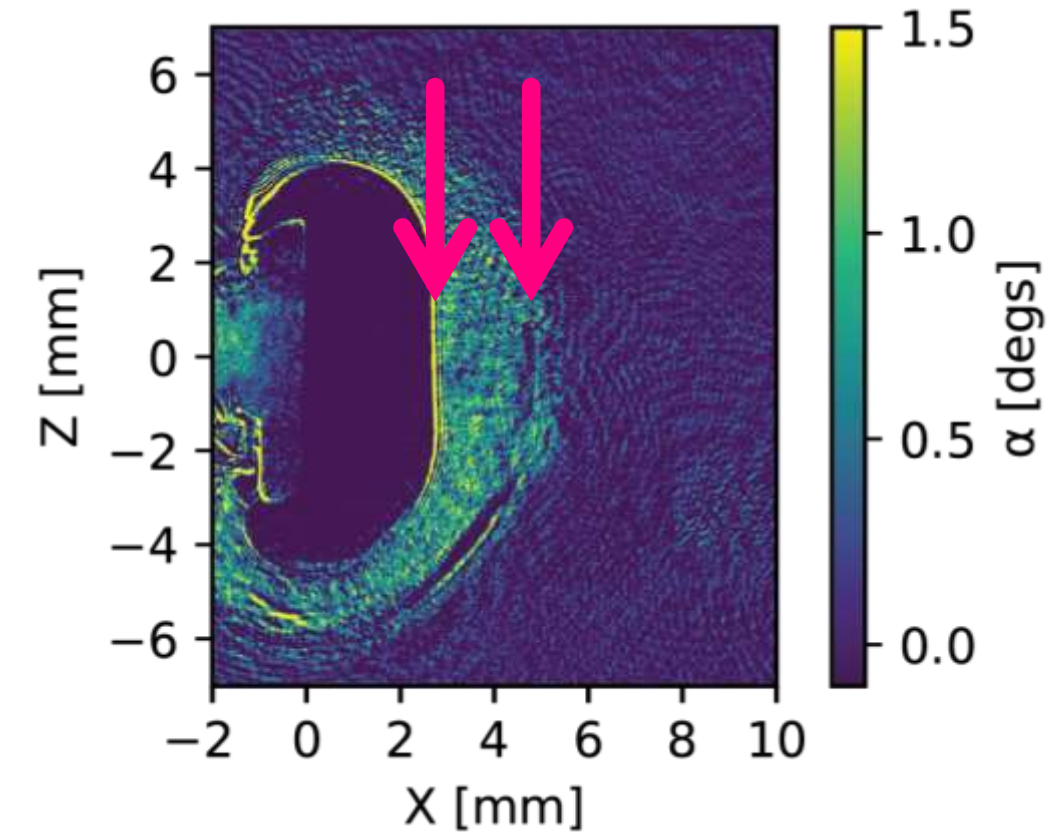
$$n_e L = \int n_e dy$$

- Combine data to back-out **weighted average** magnetic field:

$$\overline{B_y} = \frac{\alpha}{n_e L} \propto \frac{\lambda^2 \int n_e \vec{B} \cdot d\vec{y}}{\int n_e dy}$$



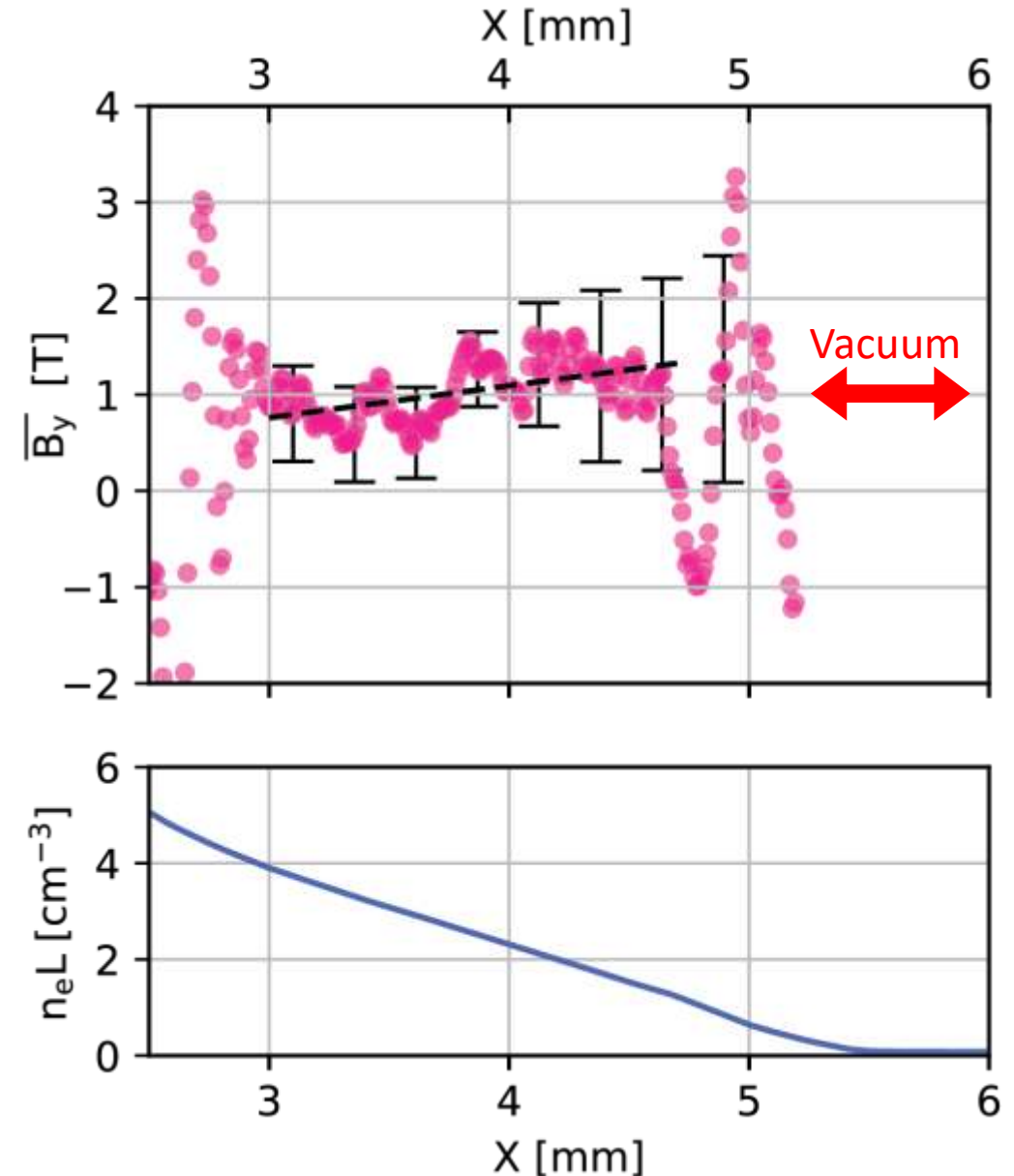
Arrows indicate caustics  $\Rightarrow$  B field not representative

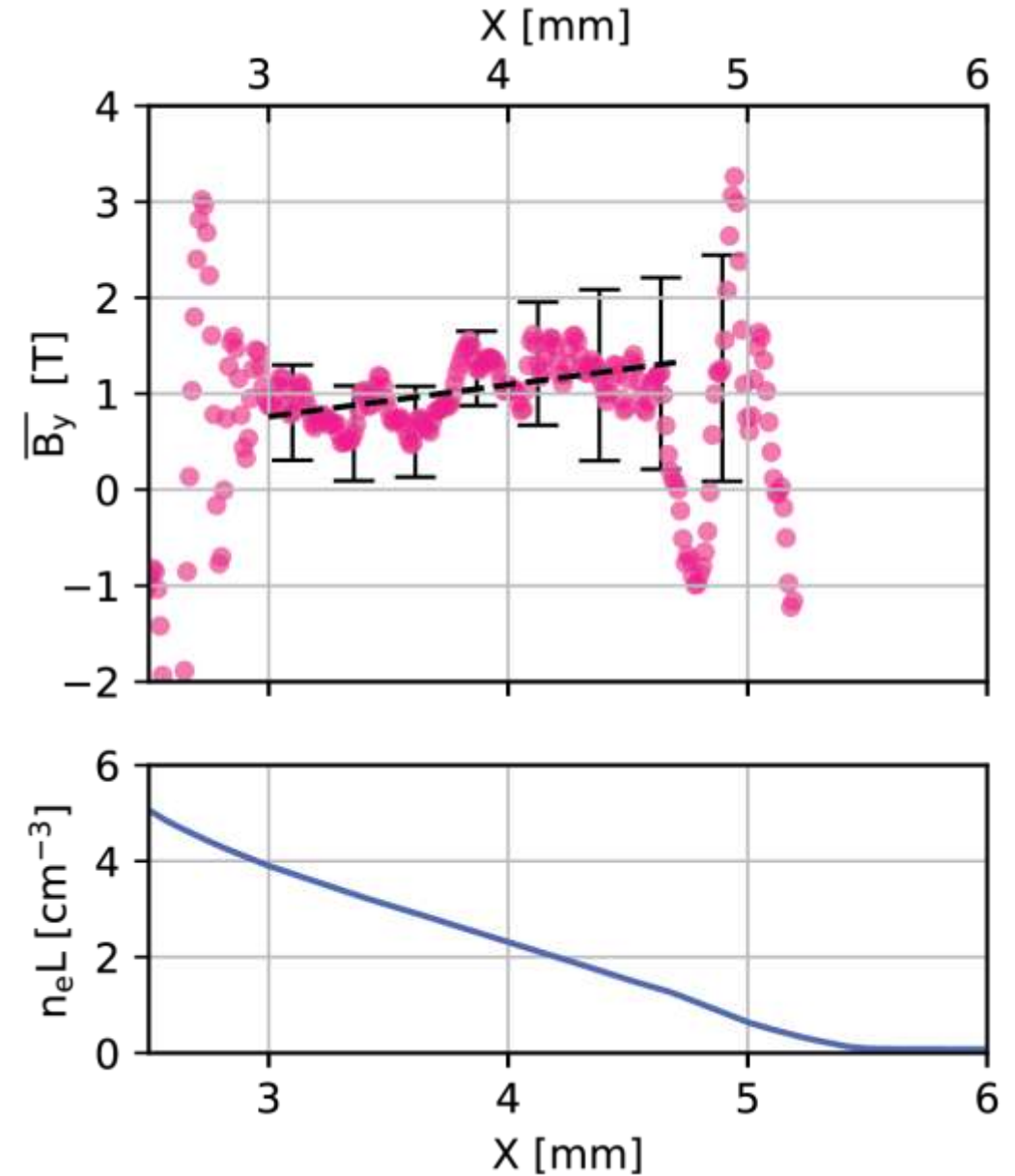
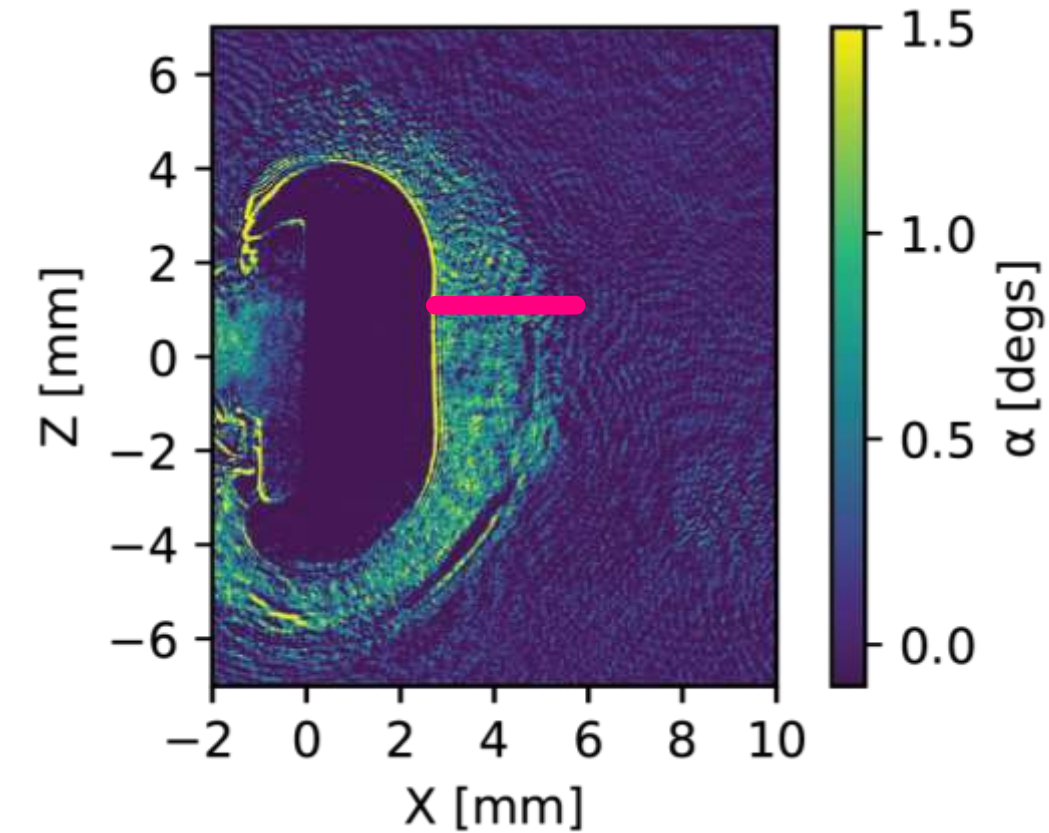


- Diagnostic measures weighted average magnetic field:

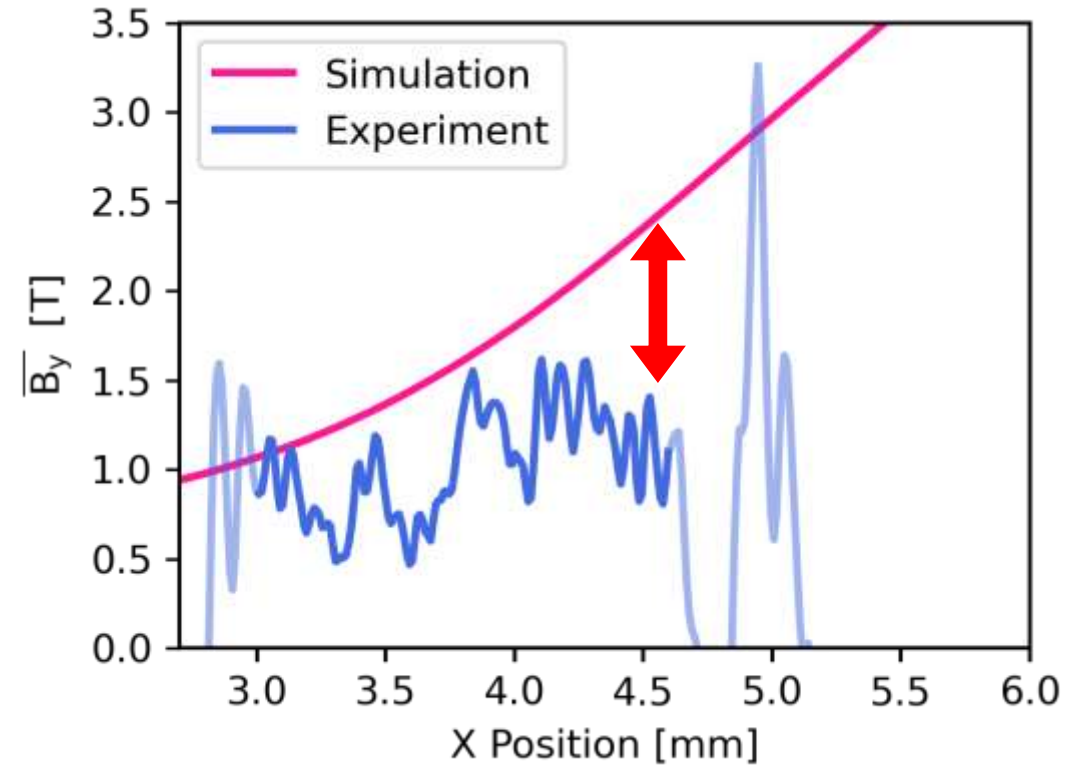
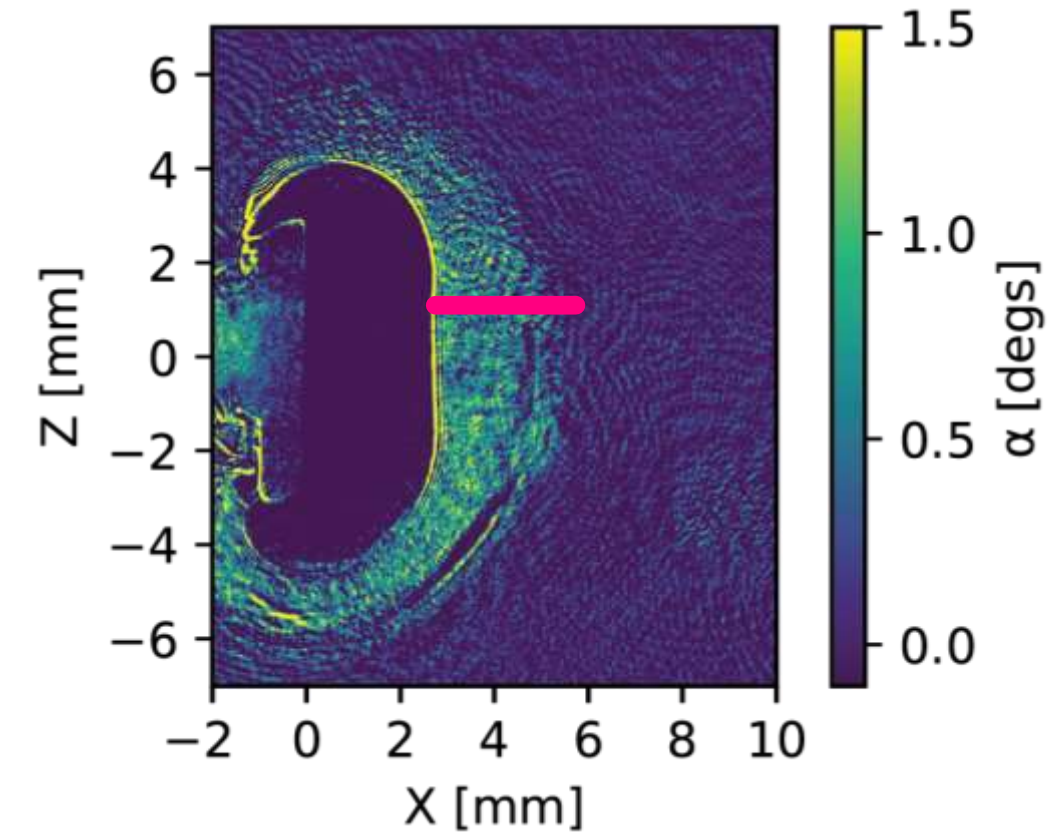
$$\overline{B_y} = \frac{\alpha}{n_e L} \propto \frac{\lambda^2 \int n_e \vec{B} \cdot d\vec{y}}{\int n_e dy}$$

- Cannot diagnose field in the vacuum ( $n_e = 0$ )
- Within region which can be probed, the field is approximately constant ( $\sim 1$  T)

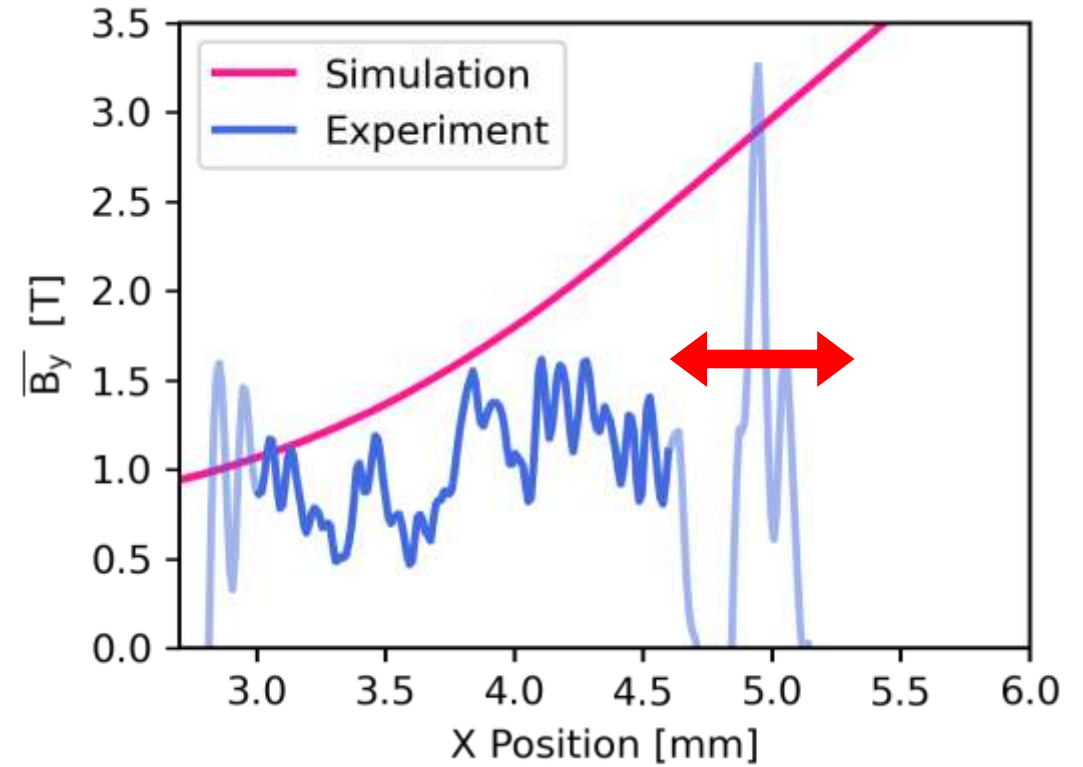
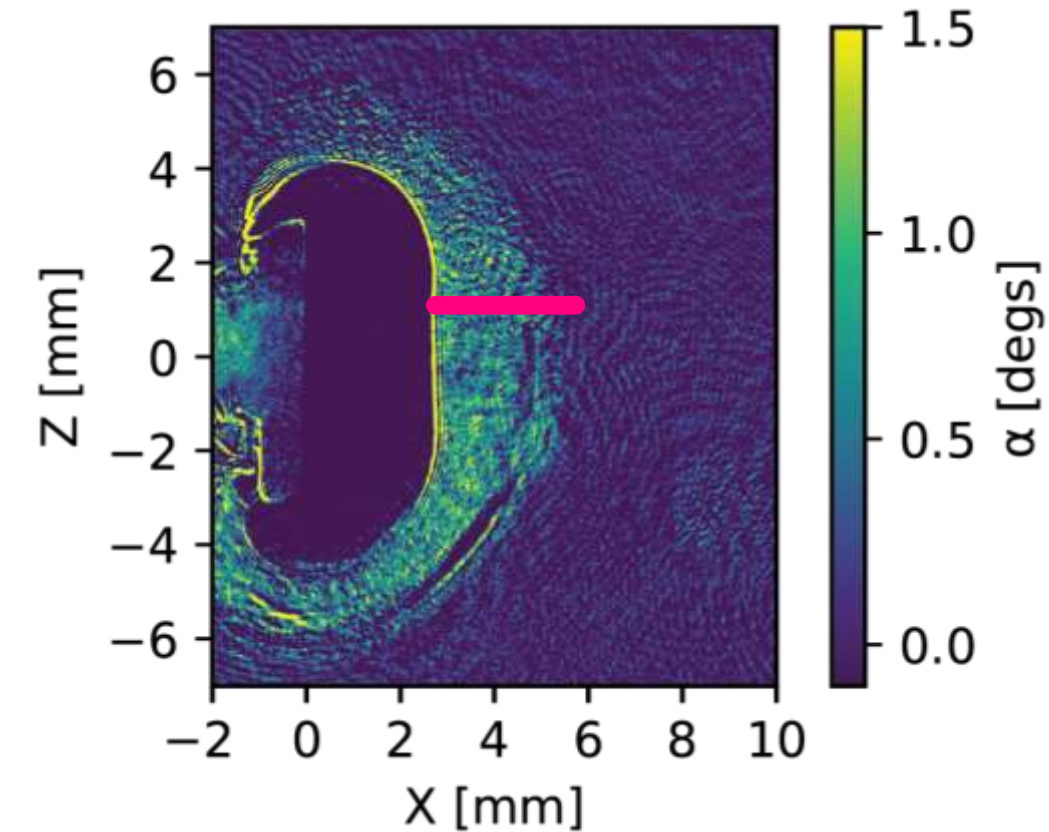








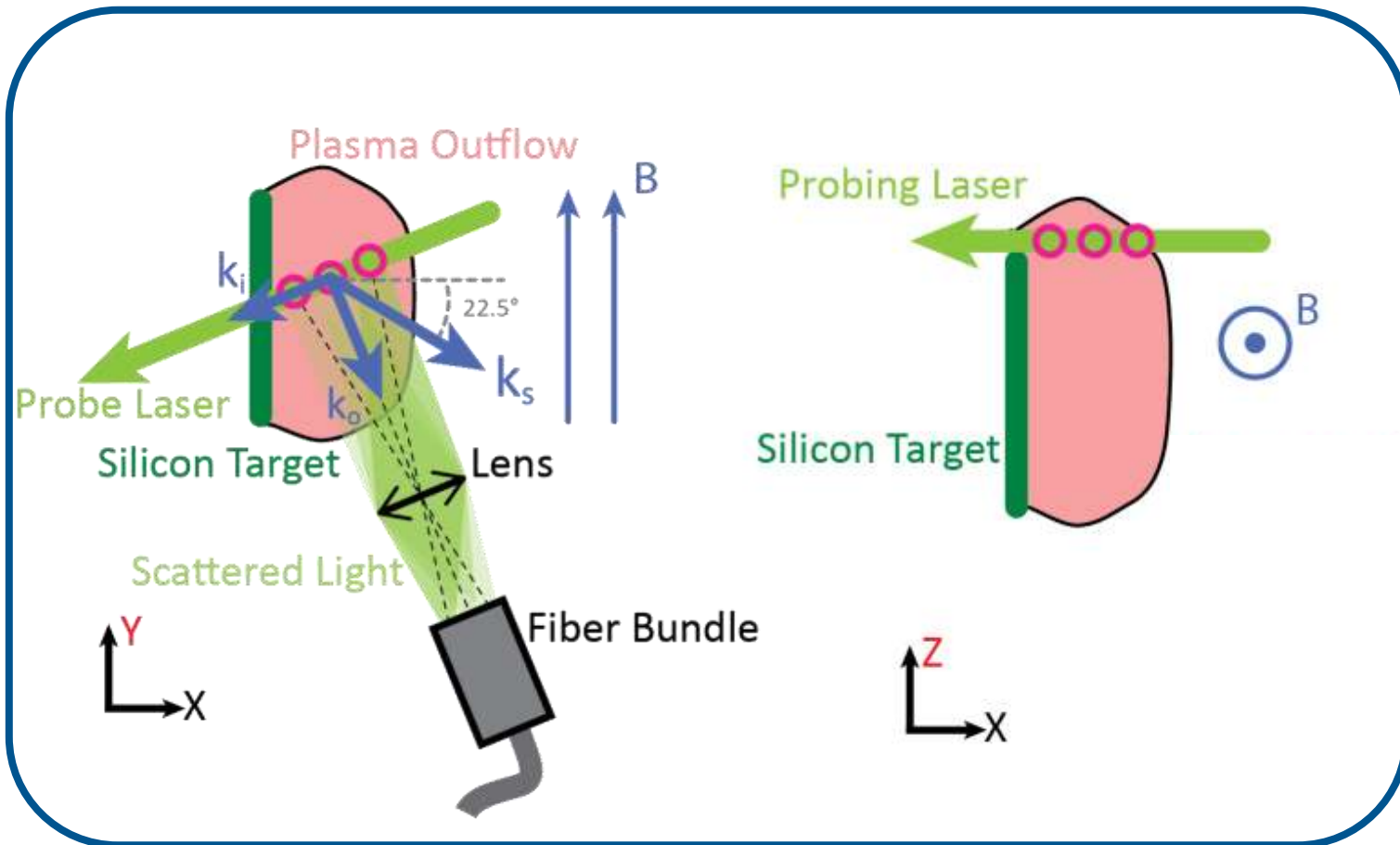
Simulated profile is more diffusive



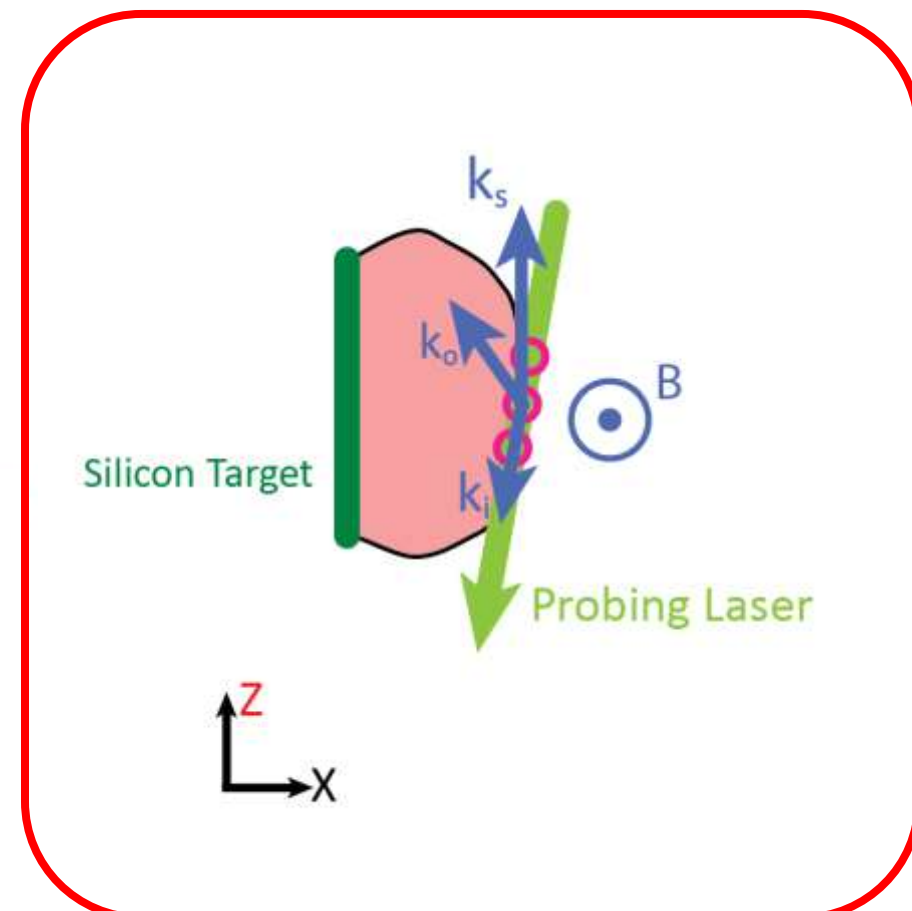
Need to diagnose vacuum boundary!

# Future work – Local Current Density Measurement

Existing setup  $\Rightarrow k_s \perp J$



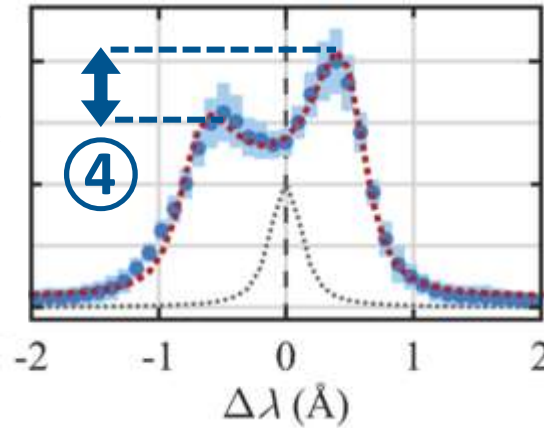
New setup  $\Rightarrow k_s \parallel J$



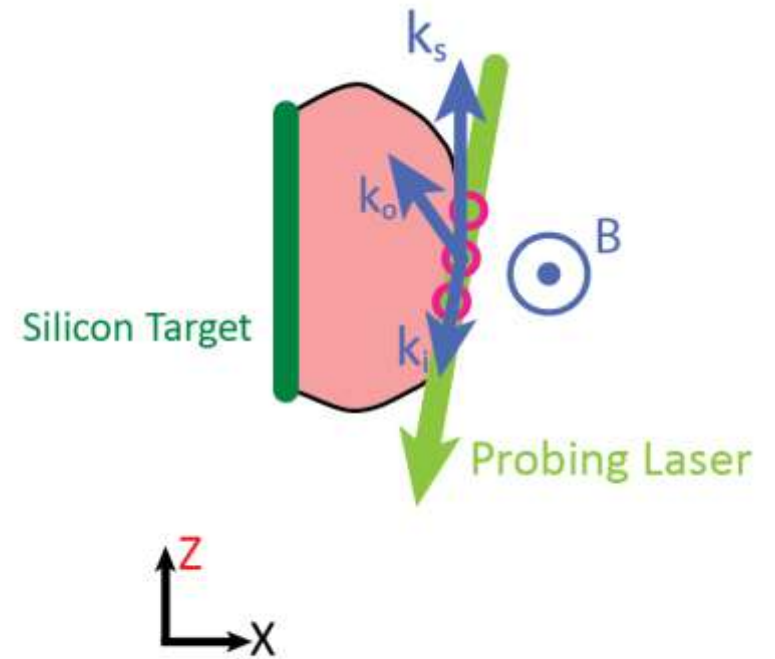
# Future work – Local Current Density Measurement

④: Peak **asymmetry** depends on  $\widehat{k}_s \cdot V_d$

L. G. Suttle *et al* "Collective Thomson scattering in pulsed-power driven HED experiments" RSI 2021.

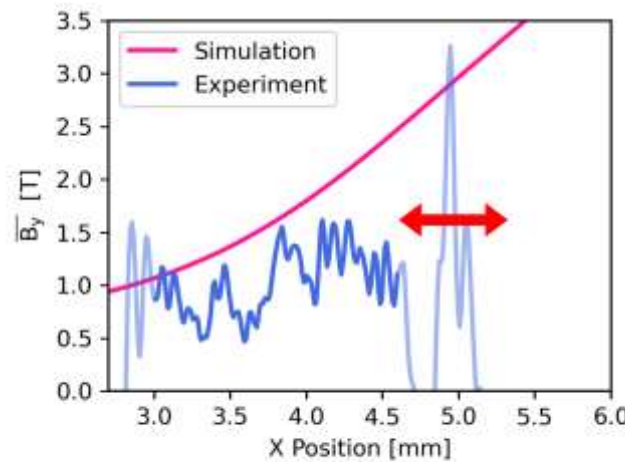


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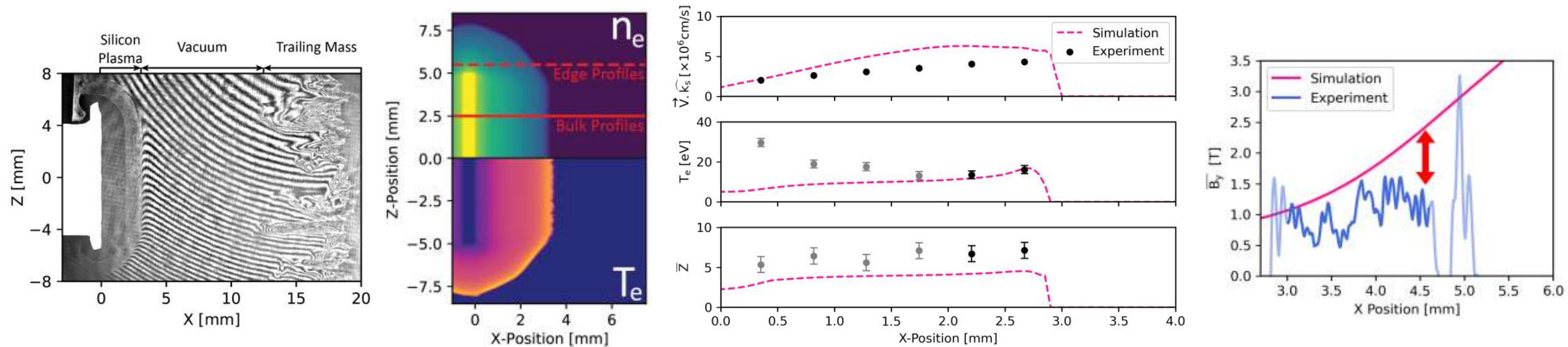
**Diagnose current at vacuum boundary with Thomson:**

- Can probe smaller  $n_e$
- Reduce  $\lambda$  for less diffraction

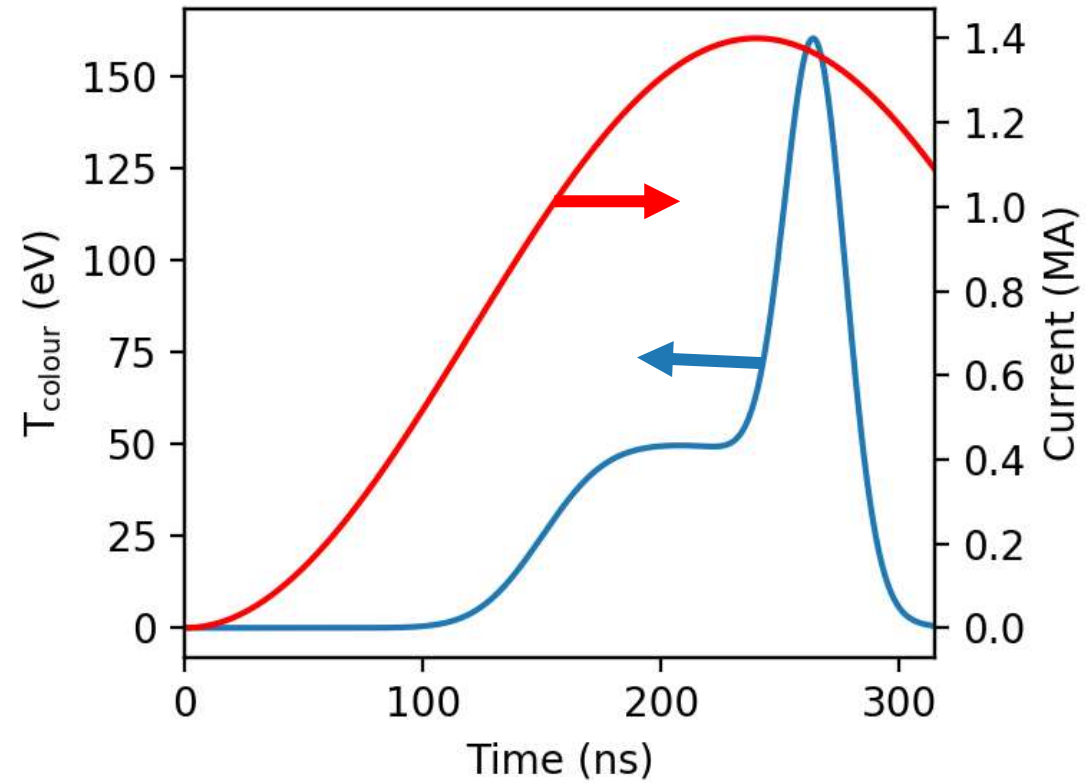


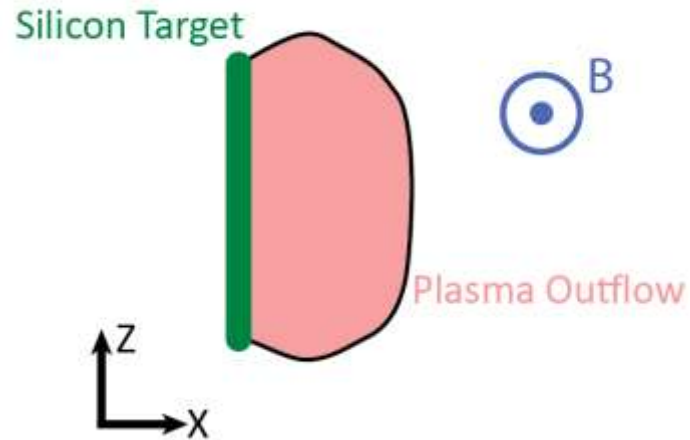
Need to diagnose vacuum boundary!

# Summary



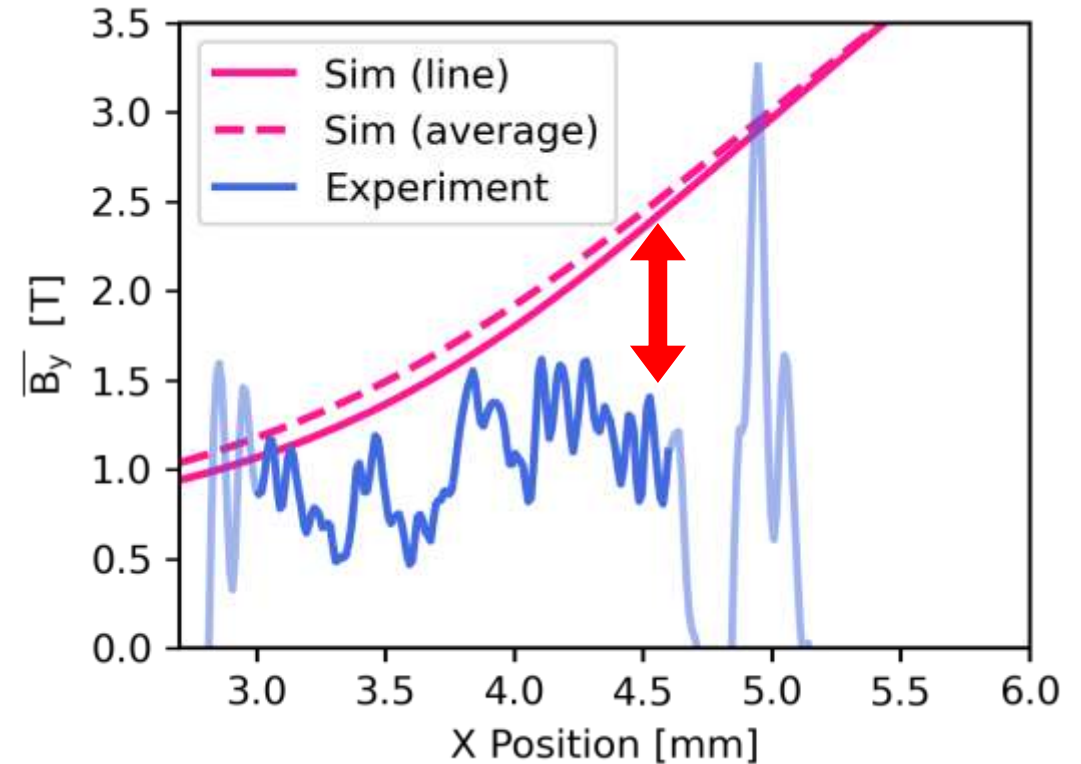
- Experimental morphology well reproduced by simulations
- Saw influence of Thomson probe heating
- Radiation field plays a role in charge state distribution (?)
- Simulated B field is more diffusive than experiment





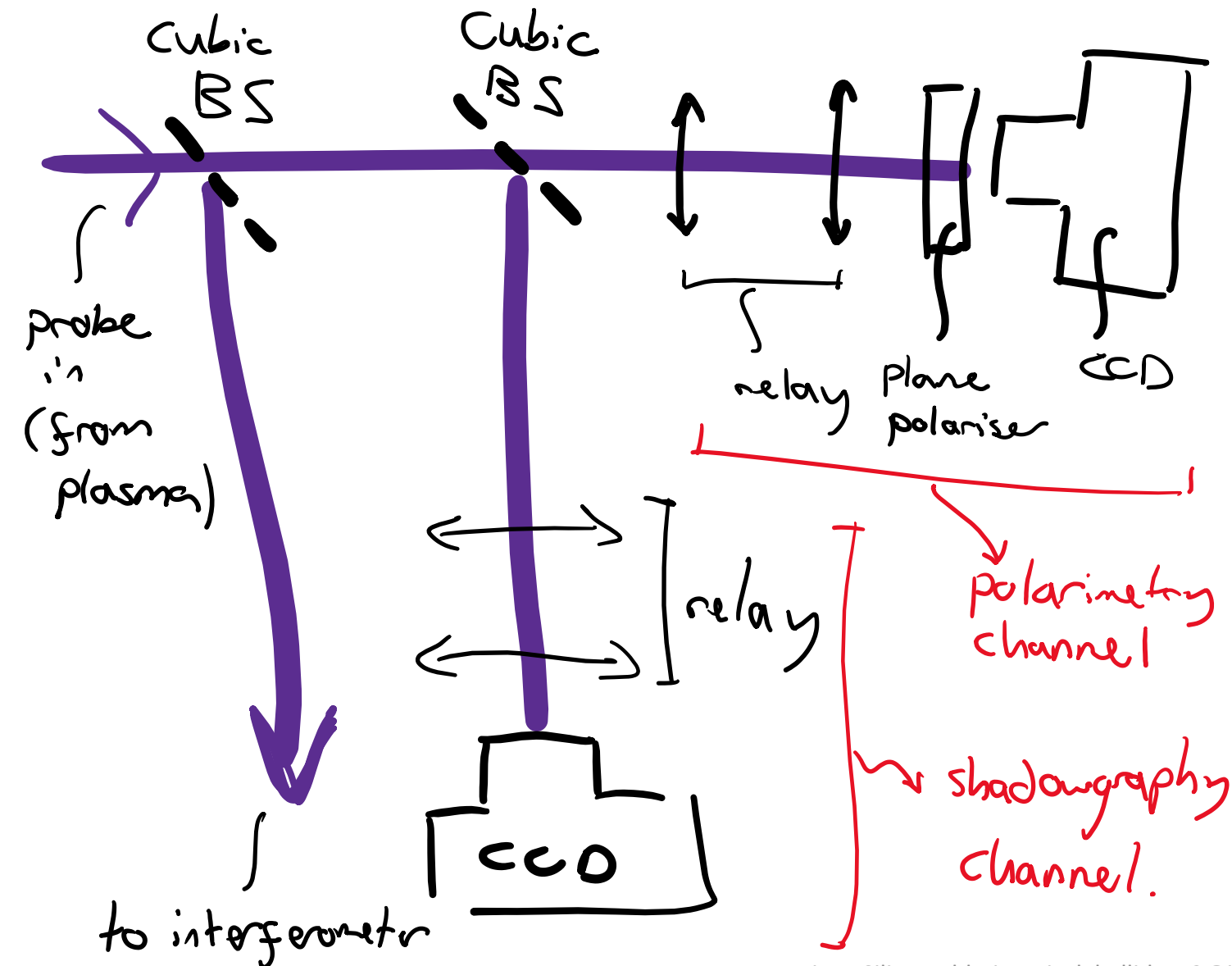
$$\overline{B_y} = \frac{\alpha}{n_e L} \propto \frac{\lambda^2 \int n_e \vec{B} \cdot d\vec{y}}{\int n_e dy}$$

$$B_{y,av} = \frac{\int n_e B_y dz}{\int n_e dz}$$



Simulated profile is more diffusive

# Faraday rotation imaging [weighted average of $B_y$ ]



In the absence of shadowgraphy effects, during the experiment, the signal measured at position  $x, z$  in the shadowgram is given by

$$I_s(x, z) = R_s(x, z)T(x, z)I_0(x, z), \quad (5)$$

where  $R_s$  is the response of the shadowgraphy camera;  $T$  is the fraction of laser light which is transmitted through the plasma; and  $I_0$  is the initial intensity of the probe. For an image taken with the same optical setup but *in the absence of plasma*, the signal measured is given by

$$I_s^*(x, z) = R_s(x, z)I_0^*(x, z). \quad (6)$$

For the polarogram, the signal during the experiment is given by

$$I_p(x, z) = R_p(x, z)T(x, z) \sin^2[\beta - \alpha(x, z)]I_0(x, z), \quad (7)$$

and the signal in the absence of plasma is

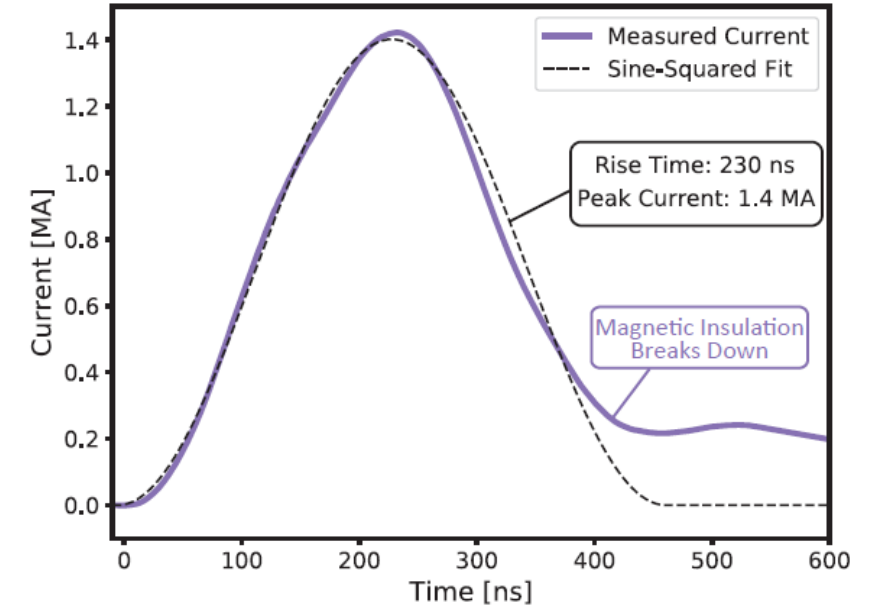
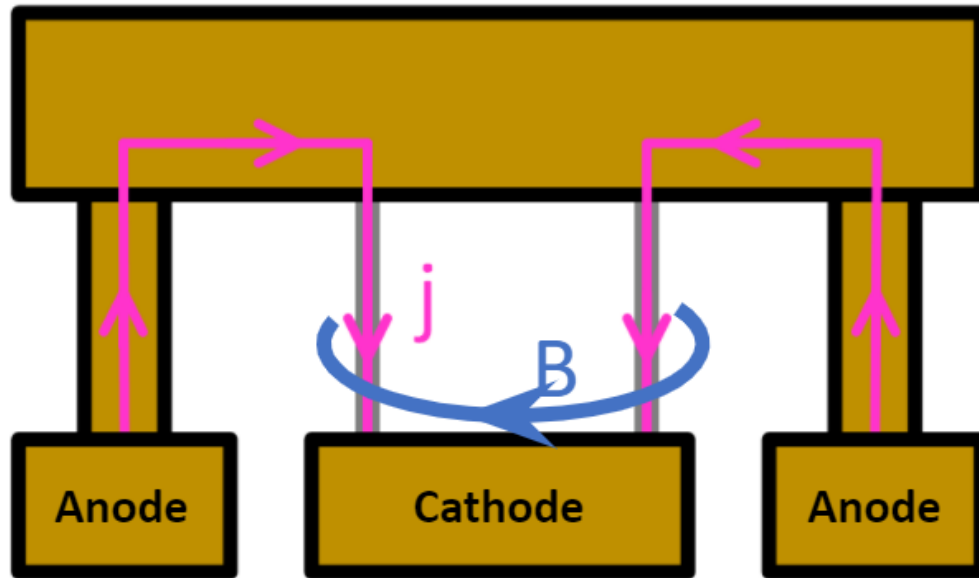
$$I_p^*(x, z) = R_p(x, z) \sin^2(\beta)I_0^*(x, z). \quad (8)$$

Combining these four equations, and solving for  $\alpha$ , yields

$$\alpha(x, z) = \beta - \arcsin \left( \sqrt{\frac{I_p(x, z)I_s^*(x, z)}{I_p^*(x, z)I_s(x, z)}} \sin(\beta) \right). \quad (9)$$



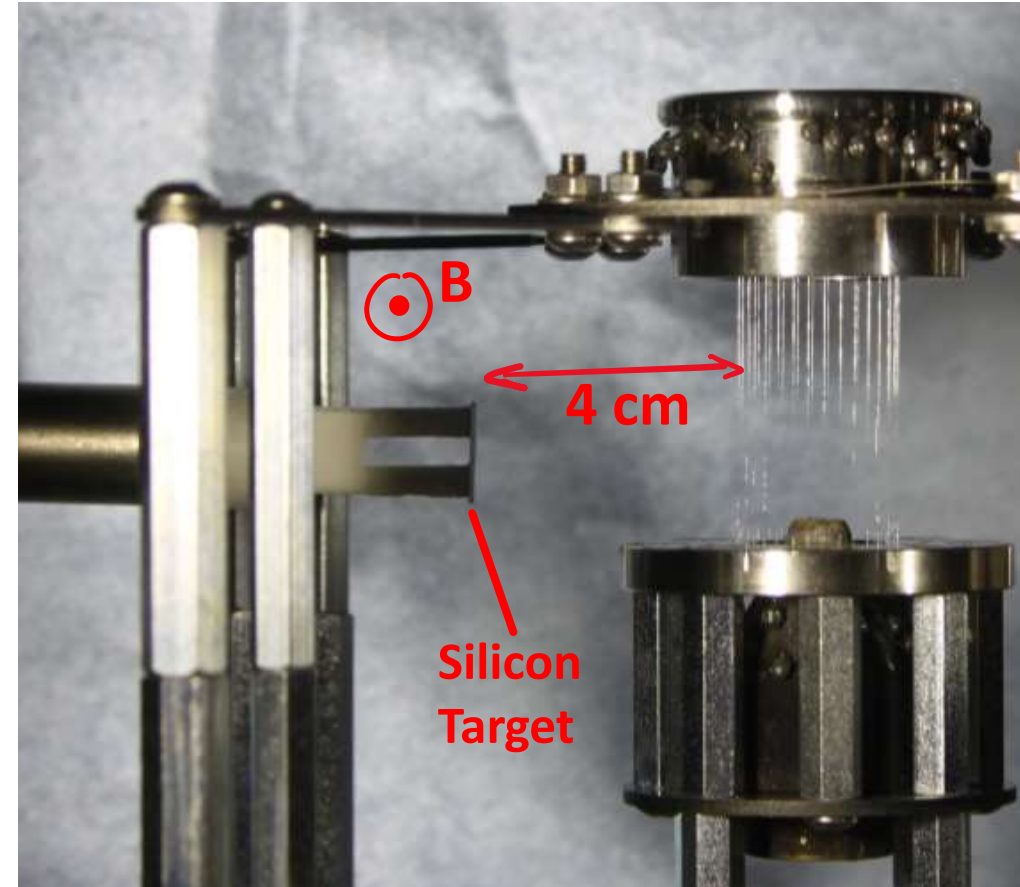
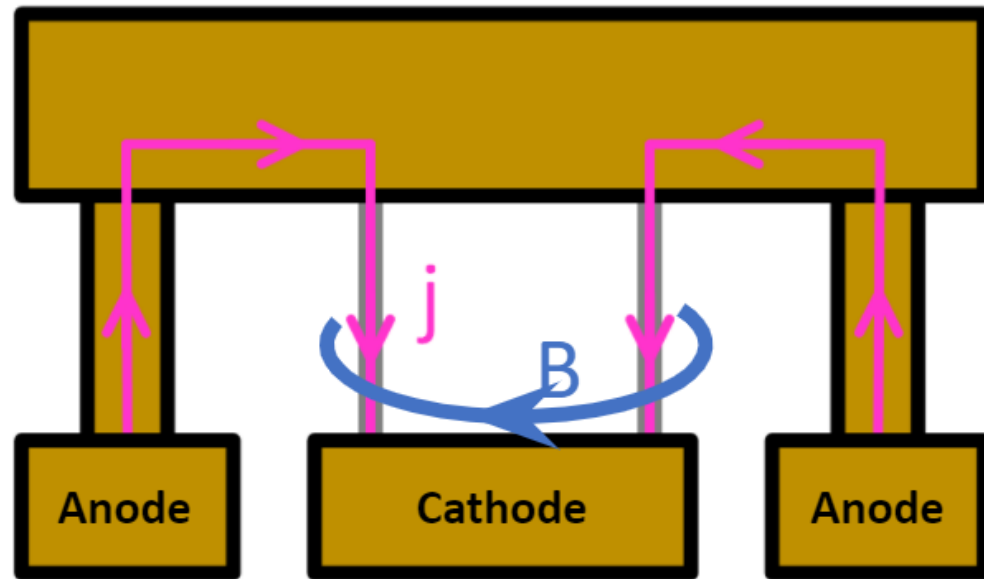
# Imploding Wire Array Z-Pinches

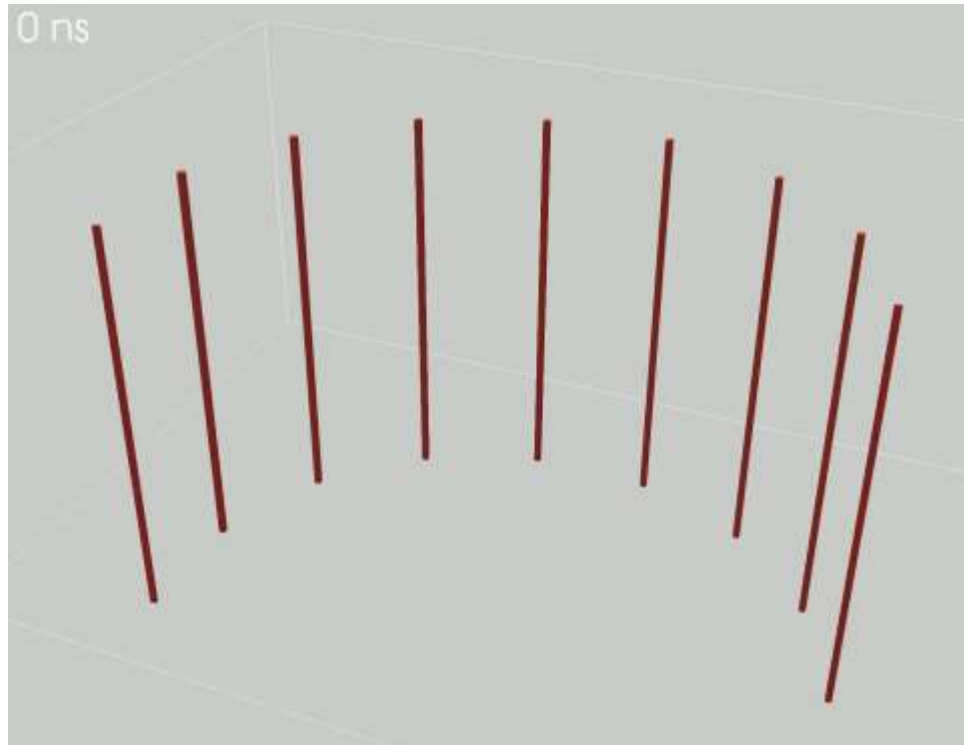


1.4 MA, 240 ns Current Pulse

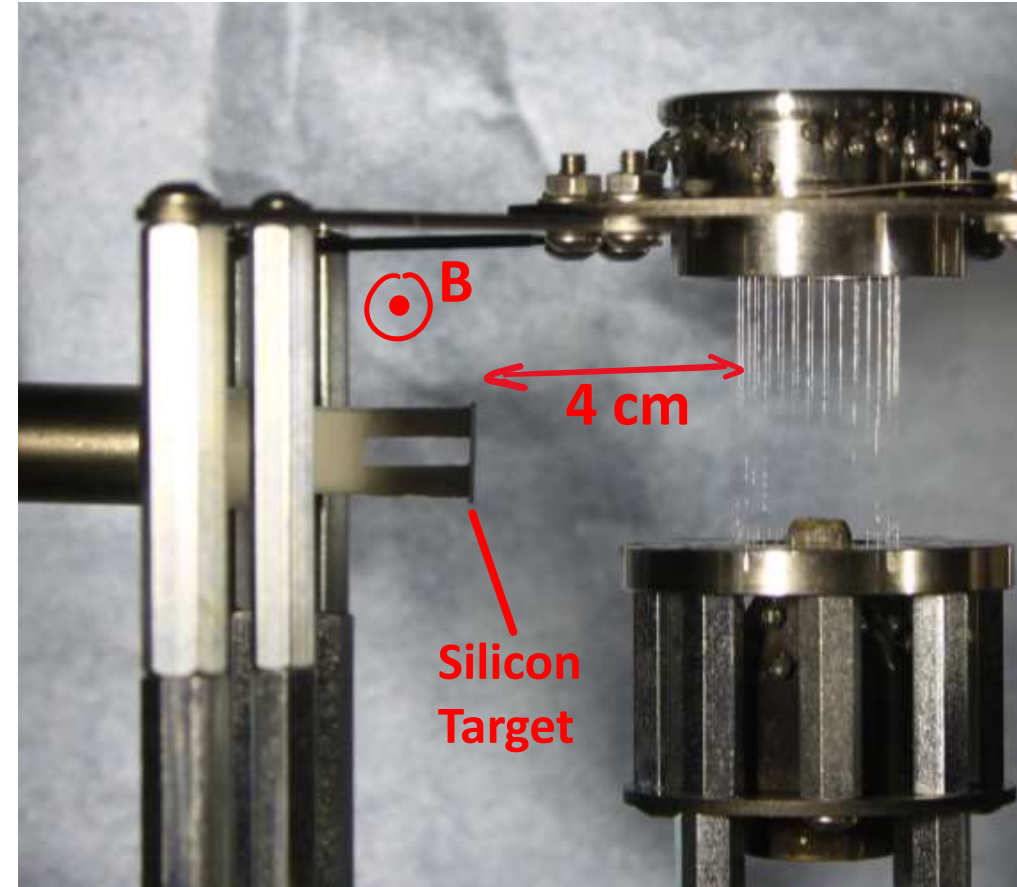
X-Ray Pulse  $\sim 1$  TW

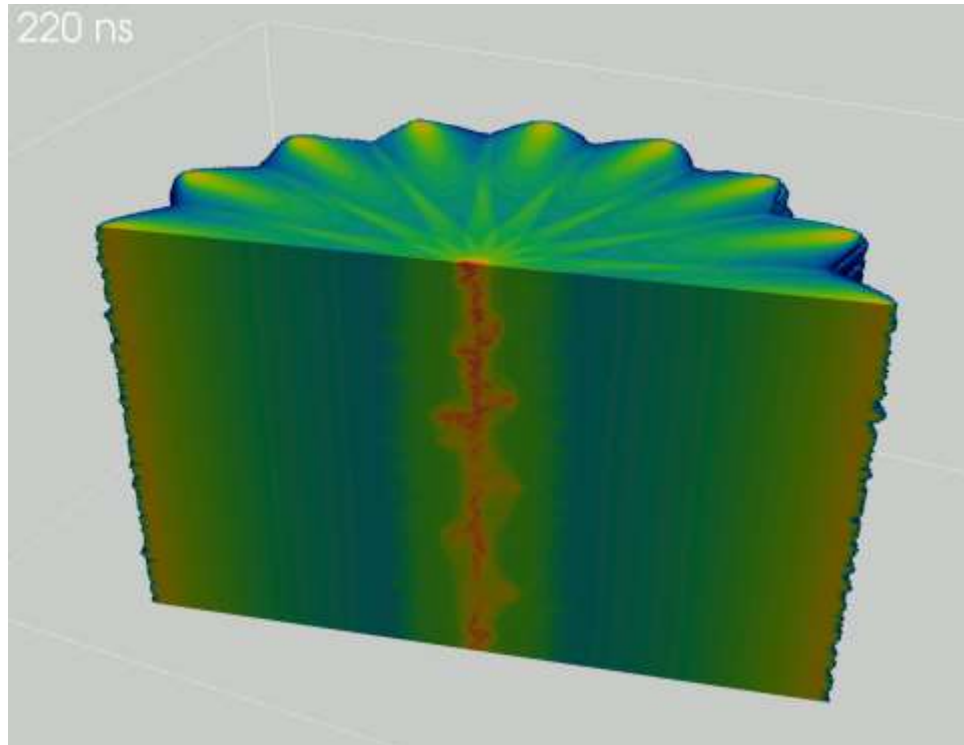
# Imploding Wire Array Z-Pinches





Mass density from Gorgon (MHD) simulation





Mass density from Gorgon (MHD)  
simulation

